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published in

Physical Review A
2019

DOI (link to publisher)

[10.1103/PhysRevA.100.052503](https://doi.org/10.1103/PhysRevA.100.052503)

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citation for published version (APA)

NL-eEDM Collaboration (2019). Lifetime measurements of the A $2\Pi_{1/2}$ and A $2\Pi_{3/2}$ states in BaF. *Physical Review A*, 100(5), 1-5. Article 052503. <https://doi.org/10.1103/PhysRevA.100.052503>

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Lifetime measurements of the $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states in BaF

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(Received 6 July 2019; published 12 November 2019)

Time-resolved detection of laser-induced fluorescence from pulsed excitation of electronic states in barium monofluoride (BaF) molecules has been performed in order to determine the lifetimes of the $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states. The method permits control over experimental parameters such that systematic biases in the interpretation of the data can be controlled to below 10^{-3} relative accuracy. The statistically limited values for the lifetimes of the $A^2\Pi_{1/2}(v=0)$ and $A^2\Pi_{3/2}(v=0)$ states are 57.1(3) and 47.9(7) ns, respectively. Although discrepancies between presently and previously measured lifetimes and theoretical values persist, the ratio of these values is in good agreement with scaling for the different excitation energies. The results serve as a benchmark for current molecular structure calculations. The investigated molecular states are of relevance for an experimental search for a permanent electric dipole moment of the electron in BaF.

DOI: [10.1103/PhysRevA.100.052503](https://doi.org/10.1103/PhysRevA.100.052503)

I. INTRODUCTION

Molecular systems, including barium monofluoride (BaF), have received significant attention in the context of experimental investigations of nuclear-spin-dependent parity-violation effects [1] and searches for electric dipole moments (EDMs) [2–6]. For this reason, activity has commenced on molecular structure calculations of BaF [7–10]. A relevant and experimentally accessible property is the radiative lifetime of electronically excited states. The measurement of lifetimes in cesium [11,12], for example, in combination with atomic structure calculations, has impacted the interpretation of atomic-parity violation [13]. We present a study of lifetimes in BaF, confirming a significant discrepancy between previous experimental results [14,15] and a recent calculation [10].

Previous measurements on radiative lifetimes in the BaF molecule have been performed in a resistance furnace using pulsed dye laser light [14–16]. These measurements were limited by contributions from collisions and background which are not related to the time-dependent fluorescence from molecules. We present a method to determine lifetimes using improved control over pulse length and timing accuracy made possible by current technology, combined with an intense and cold supersonic molecular beam.

We demonstrate a nearly-background-free method for the determination of excited-state radiative lifetimes, which are sensitive to the entire spatial extent of the wave functions. In particular, we investigate the $A^2\Pi_{1/2}(v=0)$ and $A^2\Pi_{3/2}(v=0)$ states in BaF, which govern the possibility of laser cooling of this molecular system [7,8,10,17]. Our results underline the significance of the observed discrepancy between state-of-the-

art calculations and experimentally measured lifetimes and will thus serve as a benchmark for future molecular structure calculations. In addition to the excited states of BaF, this also holds for SrF and CaF molecules, which are prototype species for laser cooling of molecules [18–21], where similar discrepancies between experimental and calculated lifetimes were found [10].

The measurement is performed on a fast internally cold supersonic molecular beam in vacuum conditions. The molecules are excited by short pulses of resonant laser radiation. Laser-induced fluorescence (LIF) is detected by a photomultiplier tube (PMT) with subnanosecond time resolution with respect to the excitation pulse. The accuracy of the timing for the laser pulses and the detection is derived from a GPS-stabilized reference clock. A detailed study of systematic biases is possible in our experiment due to a low rate of background photons in the detection systems and due to simultaneous measurement of the PMT response. The analysis shows that these biases can be controlled to a level of 10^{-3} on the yielded lifetimes.

II. METHODS

Metal-fluoride molecules, such as BaF, can be produced efficiently in a supersonic beam source [22–25]. The molecules are excited to the higher electronic state by short laser pulses (see Fig. 1). The mean lifetime value of the molecules in the excited state is extracted from the fluorescence decay after the laser pulses.

The supersonic ablation source operates at a repetition rate of 10 Hz. An Even-Lavie valve [26] creates short ($\sim 35 \mu\text{s}$) pulses of carrier gas, which is a mixture of 98% argon (Ar) and 2% sulfur hexafluoride (SF_6). The disk-shaped barium metal target is mounted inside the vacuum chamber at a distance of

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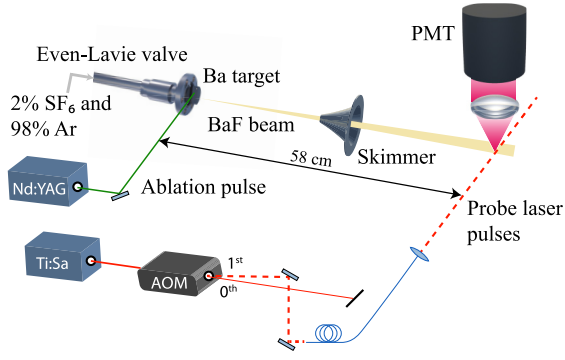


FIG. 1. Schematic overview of the experimental setup. BaF molecules are produced in a supersonic source and pass through a skimmer into the interaction region. Pulses are generated from the cw light of a Ti:sapphire (Ti:Sa) laser using the first-order diffracted beam from an acousto-optic modulator (AOM). The light from the Ti:Sa laser is delivered to the molecular beam via a 50 m single-mode optical fiber.

6 mm perpendicular to the supersonic beam. Barium atoms are ablated by 10-ns-long pulses from a Nd:YAG laser, with an energy of 20 mJ at a wavelength of 532 nm. The ablation products are entrained in the carrier gas and react with SF₆ to form BaF molecules which cool in the supersonic expansion. The molecular beam then passes through a skimmer of diameter 4 mm, located 28 cm downstream of the valve, and enters a second chamber where the molecules are detected by LIF.

The detection laser light is generated from the continuous-wave (cw) output of a Ti:sapphire laser. Pulses are created from this light by selecting the first-order diffracted beam from an acousto-optic modulator (AOM, Brimrose TEM-200-50) powered by radio frequency (rf) at 200 MHz. The pulse structure and amplitude are controlled by switching the rf (Mini-Circuits ZASWA-2-50-DR+) to obtain pulses which are synchronized to interact with the passing molecular pulse. The fraction of the first-order diffracted light beam in the fundamental Gaussian mode is coupled into a 50-m-long single-mode optical fiber, before sending it to the supersonic beam setup.

The laser pulses intersect with the molecular beam perpendicularly, exciting molecules on the $A^2\Pi_{1/2}(v=0, J=1/2) \leftarrow X^2\Sigma^+(v=0)$ (860 nm) or $A^2\Pi_{3/2}(v=0, J=3/2) \leftarrow X^2\Sigma^+(v=0)$ (815 nm) electronic transition. The fluorescence emitted is collected by a 2-inch-diameter achromatic lens with a focal length of 75 mm, which focuses the light onto an infrared-sensitive PMT (H7422-50 Hamamatsu) with an active area of 5 mm diameter. A narrow bandpass filter with 60% transmission at the resonance wavelength selects fluorescence photons. The detection chamber is internally blackened with paint that absorbs infrared light (AZ Technology MLS-85-SB) in order to reduce the background count rate. The cooled PMT has a count rate of about 30 s⁻¹ from thermionic emission and of 2000 s⁻¹ for cw laser light of power 1 mW passing through the interaction zone.

A pulse train is generated with a period of 2 μ s, a duty cycle of 2%, and a pulse length of 40 ns. Since the duration of the excitation pulse is comparable to the lifetime of the excited state, the Fourier-limited bandwidth of the pulses is comparable to the linewidth of the employed transitions.

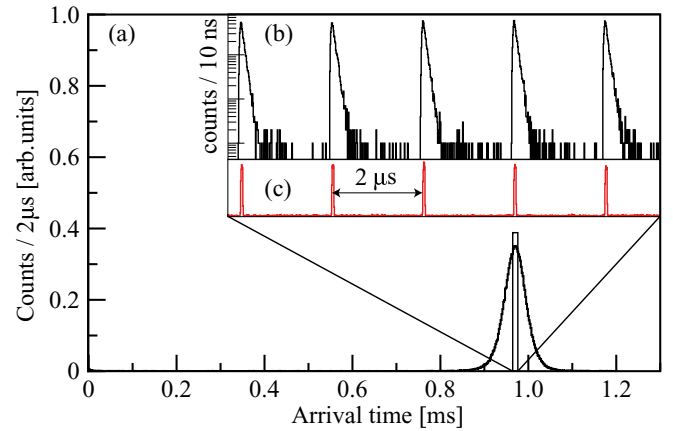


FIG. 2. Time-of-flight profile showing the accumulated fluorescence of the BaF molecular beam following excitation on the $A \leftarrow X$ transition. The histogram in panel (a) shows the arrival-time distribution of the molecular-beam fluorescence, binned at 2 μ s. Panel (b) shows the decay of the fluorescence following excitation by 40 ns laser pulses [shown in panel (c)] at repetition rate of 500 kHz, reflecting the lifetime of the $A^2\Pi_{1/2}$ state.

The 100 ms time period between two molecular pulses is split into 50 ms with the pulse train and 50 ms without laser light. The former permits the determination of the laser pulse parameters, while the latter determines the dark count rate of the PMT. The molecules arrive at the interaction region around 1 ms after the laser ablation pulse [Fig. 2(a)]. On timescales of tens of nanoseconds the response to excitation laser pulses [Fig. 2(c)] and the molecular fluorescence [Fig. 2(b)] is observed. The arrival time of each photon is recorded by a multihit-capable time-to-digital converter (Signadyne, SD-TDC-H3344-PXIe). Accurate timing synchronization is achieved by referencing all timing units to a GPS-disciplined 10 MHz rubidium frequency standard ($\Delta\nu = 10^{-12}$ in 2000 s, FS725 Stanford Research Systems).

The molecular beam has a velocity of (594 ± 18) ms⁻¹ corresponding to a travel of 60 μ m in 100 ns. Up to 100 photons per ablation shot are detected with cw detection laser light with a power of 2 mW and a $1/e^2$ diameter of ≤ 4 mm. The vast majority of the molecules are therefore excited in a volume small compared with the detection volume (set by the PMT active area of 5 mm). A systematic bias on the measured lifetime due to molecules traveling out of (or into) the detection volume is therefore negligible.

The timing parameters of the laser pulses are measured simultaneously with the same data-acquisition system as the molecular signal. The fall time (between 90% and 10% of the amplitude) of the laser pulses is additionally measured with a fast photodiode (2 GHz bandwidth) on an oscilloscope to be 11(1) ns. The ratio of the light intensity during the laser pulse to the intensity in the following several hundred nanoseconds, is determined to be larger than 1000 : 1.

III. RESULTS AND DISCUSSION

In single-photon counting with PMTs, each pulse has $\sim 1\%$ probability to be followed by a second pulse (afterpulse) [27–29]. These afterpulses cause the main systematic

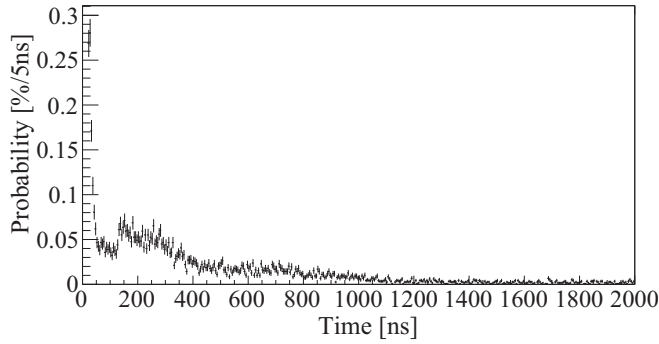


FIG. 3. Time distribution of the probability of a second photon in a time window of 2000 ns. This spectrum is extracted from the same dataset as the fluorescence data for the determination of state lifetimes.

bias in our low-background experiment. The employed multihit-capable data-acquisition system permits quantitative recording of the time distribution of these pulses (see Fig. 3) throughout all data taking periods for the lifetime measurements. The distribution of these pulses ranges over timescales of microseconds with a structure comparable to the LIF signal. This contribution, although small, is taken into account in the determination of the excited-state lifetimes.

The time distribution $N(t)$ of the number of photons recorded at time t after the trailing edge of the excitation laser pulse is described by

$$N(t) = \frac{A}{\tau} e^{-t/\tau} + B(t) + C, \quad (1)$$

where A is the number of photons recorded in the exponential decay, τ is the lifetime of the excited state, $B(t)$ is the convolution of the distribution of afterpulses (Fig. 3) with the distribution $N(t)$, and C is a constant offset due to background counts. The uncertainties on the counts per channel arise from Poissonian statistics.

A maximum likelihood fitting procedure of the function given in Eq. 1 to the experimental data yields a value of $\tau_{\Pi_{1/2}} = 57.1(3)$ ns for the lifetime (Fig. 4). The contribution

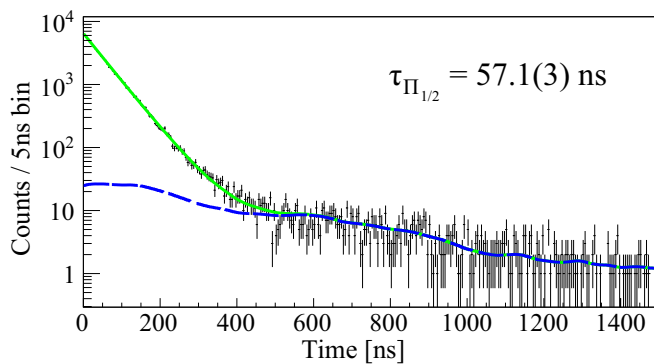


FIG. 4. Accumulated spectrum for the time distribution $N(t)$ for the $A^2\Pi_{1/2}(v=0)$ state. A maximum-likelihood fit of Eq. (1) to the data is shown by the solid (green) line. The dashed (blue) line indicates the contribution of the time-dependent background $B(t)$. A value for the lifetime of $\tau_{\Pi_{1/2}} = 57.1(3)$ ns is obtained.

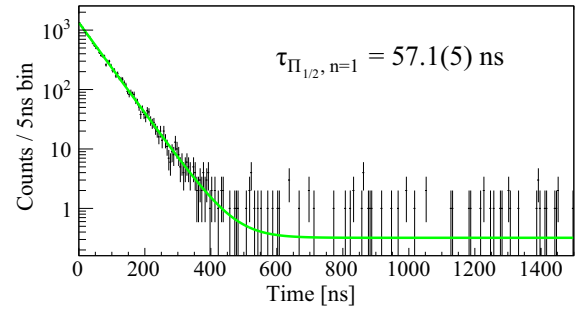


FIG. 5. Accumulated spectrum of the time distribution $N(t)$ for the $A^2\Pi_{1/2}$ state, selecting events with one photon per molecular pulse. The yielded lifetime $\tau_{\Pi_{1/2}, n=1} = 57.1(5)$ ns.

of the time-dependent term $B(t)$ is 1.6% of the total statistics. Not taking $B(t)$ into account leads to a bias of 1.6 ns on the lifetime, which is five times larger than the statistical uncertainty.

On average, three photons per molecular pulse ($\bar{n} = 3$) were recorded. For a quarter of these molecular pulses, exactly 1 photon ($n = 1$) was recorded. In this subset of the data, there is no contribution from afterpulses. The resulting spectrum is therefore analyzed with the model function (1) with the contribution $B(t)$ set to zero (see Fig. 5). For the $A^2\Pi_{1/2}$ state, the lifetime $\tau_{\Pi_{1/2}, n=1} = 57.1(5)$ ns. This number is in excellent agreement with the result obtained with the full dataset. The same analysis method applied to the recorded time spectra for the $A^2\Pi_{3/2}(v=0)$ state yields the lifetime $\tau_{\Pi_{3/2}, n=1} = 47.9(6)$ ns (see Fig. 6).

The magnetic field in the measurement region has been determined to be $<50 \mu\text{T}$. It is directed vertically, coinciding with the observation line of the PMT. The excitation laser light was elliptically polarized. In such a geometry, modulation of the fluorescence from the excited state due to Zeeman quantum beats [11] is strongly suppressed. The estimated periods of the quantum beats are 0.9 and 3.3 μs for the $A^2\Pi_{3/2}$ and the $A^2\Pi_{1/2}$ state, respectively.

Systematic bias on the measurement due to the excitation light pulse parameters was checked by variation of the duty cycle between 1% and 20% and of the repetition period between 1 and 10 μs . From a systematic scan of the start time for the fitting procedure in Fig. 4, we find that the variation of the lifetime value is one order of magnitude smaller than the

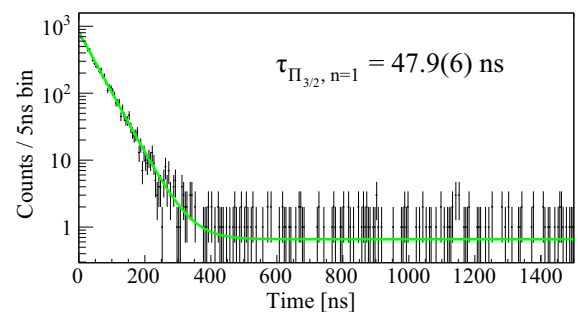


FIG. 6. Accumulated spectrum of the time distribution $N(t)$ for the $A^2\Pi_{3/2}$ state. Selecting events with one photon detected per molecular pulse yielded for the lifetime $\tau_{\Pi_{3/2}, n=1} = 47.9(6)$ ns.

TABLE I. Experimental and theoretical values for the lifetime of the $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states. The theoretical values were obtained in the framework of the NR-MRCI and X2C-MRCI methods.

Method	Lifetime [ns]	
	$A^2\Pi_{1/2}$	$A^2\Pi_{3/2}$
Supersonic beam (this work)	57.1(3)	47.9(6)
Resistance furnace	56.0(9) [15]	46.1(9) [14]
X2C-MRCI	40.4 [10]	34.7 [10]
NR-MRCI	37.8 [7]	

statistical uncertainty. Similar size contributions are observed for subsets of the data with different excitation probabilities due to a change in laser power or drift in the excitation laser frequency.

The measured lifetime values are listed with previous experimental work [14,15] and recent quantum chemistry calculations [7,10] in (Table I). The lifetimes of the excited states of BaF were calculated by using the multireference configuration-interaction approach in a nonrelativistic framework (NR-MRCI) [7], thus yielding the same result for the $A^2\Pi_{1/2}$ and the $A^2\Pi_{3/2}$ states. In a more recent investigation [10], we used relativistic MRCI, where the traditional four-component Dirac-Coulomb (DC) Hamiltonian was replaced by the exact two-component Hamiltonian (X2C-MRCI) [30,31]. In both works, experimental level energies were used along with calculated transition dipole moments to determine the lifetimes. The X2C-MRCI results reproduce the lifetime ratio very well, as a result of the accurate treatment of spin-orbit effects. However, the predicted lifetimes underestimate the experimental values by about 30%, due to the shortcoming of the MRCI approach in the treatment of the transition dipole moments (TDMs). The quality of the calculated TDMs is highly sensitive to the level of treatment of electron correlation, and a multireference coupled-cluster approach, such as the Fock-space coupled cluster (FSCC) [32] is expected to yield higher-accuracy results. An implementation that will allow use of FSCC for TDM calculations is presently under way [33].

The main decay channel of the investigated excited states is an electric dipole transition to the $X^2\Sigma^+(\nu = 0)$ ground state. In this case, scaling of lifetimes is expected to be proportional to the third power of the fluorescence wavelength

TABLE II. Comparison between the ratios of the lifetimes $\tau_{\Pi_{3/2}}/\tau_{\Pi_{1/2}}$ obtained in this work, previously measured experimental work [14,15], and the X2C-MRCI theoretical calculations with the expected value from λ^3 -scaling of the transition wavelengths.

Method	Lifetime ratio
This work	0.839(14)
λ^3 -scaling	0.8531
Resistance furnace [14,15]	0.823(25)
X2C-MRCI [10]	0.859

(λ^3) ($\lambda = 859.79$ nm and $\lambda = 815.45$ nm for the $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states, respectively). We compare the ratio of the lifetimes of the two excited states with the previously measured experimental work [14,15], the X2C-MRCI theoretical calculations, and the expected value from λ^3 scaling. Within the experimental accuracy, our experimental result is in agreement with this scaling (Table II).

IV. CONCLUSION

A measurement of the lifetimes of the low-lying electronically excited $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states in the BaF molecule was performed in the context of the experimental search for a permanent electric dipole moment of the electron (NL-eEDM) [6]. The statistically limited accuracies on the lifetimes of the $A^2\Pi_{1/2}$ and $A^2\Pi_{3/2}$ states are 0.5% and 1.3%, respectively. The main systematic uncertainty arises from the imperfection of photomultipliers in photon-counting mode, in particular their afterpulses. By measurement of the afterpulse spectrum of the employed PMT, the relative contributions to the measured lifetime values could be reduced to the 10^{-3} level. Methods which permit identification and control of systematic effects are crucial in precision measurements and their interpretation. Careful design of the experimental procedure and the use of state-of-the-art technology enables the execution of such measurements with low sensitivity to systematic bias.

ACKNOWLEDGMENTS

The NL-eEDM consortium receives program funding (EEDM-166) from the Netherlands Organisation for Scientific Research (NWO). We thank C. J. G. Onderwater for discussions on the statistical treatment of the data.

- [1] E. Altuntaş, J. Ammon, S. B. Cahn, and D. DeMille, *Phys. Rev. A* **97**, 042101 (2018).
- [2] J. J. Hudson, D. M. Kara, I. J. Smallman, B. E. Sauer, M. R. Tarbutt, and E. A. Hinds, *Nature (London)* **473**, 493 (2011).
- [3] W. B. Cairncross, D. N. Gresh, M. Grau, K. C. Cossel, T. S. Roussy, Y. Ni, Y. Zhou, J. Ye, and E. A. Cornell, *Phys. Rev. Lett.* **119**, 153001 (2017).
- [4] A. C. Vutha, M. Horbatsch, and E. A. Hessels, *Phys. Rev. A* **98**, 032513 (2018).
- [5] The ACME Collaboration, V. Andreev, D. G. Ang, D. DeMille, J. M. Doyle, G. Gabrielse, J. Haefner, N. R. Hutzler, Z. Lasner,

- C. Meisenhelder, B. R. O'Leary, C. D. Panda, A. D. West, E. P. West, and X. Wu, *Nature (London)* **562**, 355 (2018).
- [6] The NL-eEDM collaboration, P. Aggarwal, H. L. Bethlem, A. Borschevsky, M. Denis, K. Esajas, P. A. B. Haase, Y. Hao, S. Hoekstra, K. Jungmann, T. B. Meijknecht, M. C. Mooij, R. G. E. Timmermans, W. Ubachs, L. Willmann, and A. Zapara, *Eur. Phys. J. D* **72**, 197 (2018).
- [7] S. Kang, F. Kuang, G. Jiang, and J. Du, *Mol. Phys.* **114**, 810 (2016).
- [8] T. Chen, W. Bu, and B. Yan, *Phys. Rev. A* **94**, 063415 (2016).

- [9] S. Tohme and M. Korek, *J. Quant. Spectrosc. Radiat. Transfer* **167**, 82 (2015).
- [10] Y. Hao, L. Pasteka, L. Visscher, P. Aggarwal, H. L. Bethlem, A. Boeschoten, A. Borschevsky, M. Denis, K. Esajas, S. Hoekstra, K. Jungmann, V. R. Marshall, T. B. Meijknecht, M. C. Mooij, R. G. E. Timmermans, A. Touwens, W. Ubachs, L. Willmann, and A. Zapara, *J. Chem. Phys.* **151**, 034302 (2019).
- [11] L. Young, W. T. Hill, S. J. Sibener, S. D. Price, C. E. Tanner, C. E. Wieman, and S. R. Leone, *Phys. Rev. A* **50**, 2174 (1994).
- [12] G. Toh, J. A. Jaramillo-Villegas, N. Glotzbach, J. Quirk, I. C. Stevenson, J. Choi, A. M. Weiner, and D. S. Elliott, *Phys. Rev. A* **97**, 052507 (2018).
- [13] M. S. Safronova, D. Budker, D. DeMille, D. F. J. Kimball, A. Derevianko, and C. W. Clark, *Rev. Mod. Phys.* **90**, 025008 (2018).
- [14] L. E. Berg, T. Olsson, J. C. Chanteloup, A. Hishikawa, and P. Royen, *Mol. Phys.* **79**, 721 (1993).
- [15] L. E. Berg, N. Gador, D. Husain, H. Ludwigs, and P. Royen, *Chem. Phys. Lett.* **287**, 89 (1998).
- [16] P. J. Dagdigian, H. W. Cruse, and R. N. Zare, *J. Chem. Phys.* **60**, 2330 (1974).
- [17] L. Xu, B. Wei, Y. Xia, L. Deng, and J. Yin, *Chin. Phys. B* **26**, 033702 (2017).
- [18] E. S. Shuman, J. F. Barry, and D. DeMille, *Nature (London)* **467**, 820 (2010).
- [19] M. T. Hummon, M. Yeo, B. K. Stuhl, A. Collopy, Y. Xia, and J. Ye, *Phys. Rev. Lett.* **110**, 143001 (2013).
- [20] V. Zhelyazkova, A. Cournol, T. E. Wall, A. Matsushima, J. J. Hudson, E. A. Hinds, M. R. Tarbutt, and B. E. Sauer, *Phys. Rev. A* **89**, 053416 (2014).
- [21] J. Lim, J. R. Almond, M. A. Trigatzis, J. A. Devlin, N. J. Fitch, B. E. Sauer, M. R. Tarbutt, and E. A. Hinds, *Phys. Rev. Lett.* **120**, 123201 (2018).
- [22] M. Tarbutt, J. J. Hudson, B. E. Sauer, E. A. Hinds, V. A. Ryzhov, V. L. Ryabov, and V. F. Ezhov, *J. Phys. B: At., Mol. Opt. Phys.* **35**, 5013 (2002).
- [23] T. C. Steimle, S. Frey, A. Le, D. DeMille, D. A. Rahmlow, and C. Linton, *Phys. Rev. A* **84**, 012508 (2011).
- [24] S. C. Mathavan, A. Zapara, Q. Esajas, and S. Hoekstra, *ChemPhysChem* **17**, 3709 (2016).
- [25] A. Cournol, P. Pillet, H. Lignier, and D. Comparat, *Phys. Rev. A* **97**, 031401(R) (2018).
- [26] U. Even, *Adv. Chem.* **2014**, 636042 (2014).
- [27] H. C. Burstyn, *Rev. Sci. Instrum.* **51**, 1431 (1980).
- [28] M. Zhao, L. Jin, B. Chen, Y. Ding, H. Ma, and D. Chen, *Appl. Opt.* **42**, 4031 (2003).
- [29] K. Ishii and T. Tahara, *Opt. Express* **23**, 32387 (2015).
- [30] M. Iliáš and T. Saue, *J. Chem. Phys.* **126**, 064102 (2007).
- [31] T. Saue, *ChemPhysChem* **12**, 3077 (2011).
- [32] U. Kaldor and E. Eliav, *Adv. Quantum Chem.* **31**, 313 (1998).
- [33] L. Visscher (private communication).