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LETTER TO THE EDITOR

High-resolution VUV-laser spectroscopic study of the $B^1\Sigma_u^+(v' = 0-2) \leftarrow X^1\Sigma_g^+(v'' = 0)$ Lyman bands in H_2 and HD

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Online at stacks.iop.org/JPhysB/39/L195**Abstract**

Spectroscopic measurements on the $B^1\Sigma_u^+(v' = 0-2) \leftarrow X^1\Sigma_g^+(v'' = 0)$ Lyman bands in H_2 and the $B^1\Sigma^+(v' = 0-2) \leftarrow X^1\Sigma^+(v'' = 0)$ Lyman bands in HD are presented using a narrow bandwidth vacuum ultraviolet laser system combined with an accurate frequency calibration. These measurements complete the recently published data on the Lyman frequency transitions of Philip *et al* (2004 *Can. J. Chem.* **82** 713–22).

1. Introduction

H_2 is by far the most abundant molecule in the cosmos and its characteristic electronic systems, the $B^1\Sigma_u^+ \leftarrow X^1\Sigma_g^+$ and $C^1\Pi_u \leftarrow X^1\Sigma_g^+$ band systems, known as the Lyman and Werner bands, are the most prominent spectroscopic absorption features to be observed in outer space. Because the Earth's atmosphere is not transparent in the vacuum ultraviolet (VUV) and extreme ultraviolet (XUV) ranges, it took until 1970 for the first spectral recording of H_2 in space using a rocket-borne spectrometer in the wavelength range between 100 nm and 110 nm [1]. Later a large number of lines in the VUV spectrum of molecular hydrogen was observed with the *Copernicus* satellite telescope [2]. The Hubble Space Telescope, equipped with a high-resolution XUV spectrometer, detected vibrationally excited molecular hydrogen in interstellar space, in the direction of ζ -Ophiuchi [3]. Similarly the Galileo orbiter, also carrying an extreme ultraviolet spectrometer, recorded H_2 spectra from the north and south polar regions of Jupiter [4]. Since 1999 the *Far-Ultraviolet Spectroscopic Explorer* is in orbit, detecting H_2 from nearby and remote locations in the universe [5]. In all these observations, the Lyman bands are the prominent features under investigation.

In the past decade H_2 absorption spectra of quasi-stellar objects have been recorded with ground-based telescopes; the high redshifts ($z = 2-3$) allow for registering the Lyman

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bands in the visible domain, where the light penetrates through the Earth's atmosphere. These observations are made in view of the possibility that the proton-to-electron mass ratio $\mu = M_p/m_e$ has varied on a cosmological time scale [6–11]. Such observations allow for a comparison of spectra where the absorption has taken place some 12 Gyr ago, and spectra recorded in the laboratory in the modern epoch. The most stringent constraint on a possible variation of the proton-to-electron mass ratio following from such a comparison is $\Delta\mu/\mu = (-0.5 \pm 3.6) \times 10^{-5}$ (2σ) [12]. During these evaluations the need for extremely accurate zero-Doppler transition wavelengths in the Lyman and Werner band systems has emerged.

The Ottawa group around Herzberg performed a number of high-resolution absorption and emission studies on the Lyman bands [13–16], which were followed by the elaborate emission studies of the Meudon group at somewhat higher resolution [17, 18]. Baig and Connerade [19] performed an absorption study on H₂ using synchrotron radiation, while Jungen *et al* [20] determined level energies for the B, $v = 0$ and 1 states by cascaded infrared transitions, using Fourier-transform infrared (FTIR) spectroscopy. The Amsterdam group performed extensive wavelength calibration studies on the Lyman bands of H₂ [21, 22] and HD [23], using a laser-based VUV source allowing for an absolute uncertainty of the transition frequencies of better than 0.1 cm^{-1} . Since the resolution of the Amsterdam VUV laser system was improved by a factor of 20 [24], investigations of the Lyman bands for $v' = 2$ –18 could be performed, yielding absolute accuracies of 10^{-7} or better [25]. In the latter study, problems were encountered generating wavelengths in the range from 105 nm to 110 nm. Moreover, in that range no saturated I₂ lines, to be used for reference calibration [26, 27], were available. These issues have now been resolved and here we present accurate calibrations of Lyman lines in the range from 105 nm to 110 nm, for H₂ as well as for HD.

2. Experiment

The experimental setup is described in the previous papers [24, 28] and is only summarized here. The tunable cw output of a 380 D Spectra Physics ring dye laser with Stabilok, which is operated with a 10 W Spectra Physics Millennia X Nd:YVO₄ laser as a pump laser and with DCM as a gain medium, is transmitted (via a single-mode optical fibre) into a pulsed dye amplifier, pumped by a seeded Spectra Physics Nd:YAG laser (GCR-5, repetition rate 10 Hz) and producing nearly Fourier-transform limited laser pulses with a length of about 5 ns. The pulsed amplification consists of three dye cells operated with DCM. The amplified laser beam is frequency doubled using a KD*P crystal. The UV beam (1–2 mJ per pulse) is separated into two beams by placing a metallic needle at the centre of the laser beam. The two UV beams are focused into a vacuum chamber. The VUV is produced at the focus of the UV beam in a pulsed supersonic gas expansion of xenon using a piezo-disc home-built pulsed valve (design described in [29]) with a stagnation pressure of approximately 4 bar. After frequency tripling, the VUV beam, which has a bandwidth of $\Gamma \approx 250 \text{ MHz}$ [28], is located spatially between the two fundamental UV beams. The VUV beam is led through an adjustable slit and crossed perpendicularly with a pulsed, skimmed probe gas beam originating from a supersonic gas expansion using a General valve nozzle (series 9, 0.8 mm orifice) with a stagnation pressure of approximately 1 bar. The H₂ or HD probe gases are seeded with xenon in a 1:3 mixture (H₂/HD:Xe) or used in a pure form. The distance between the probe gas nozzle and the used skimmer (opening 2 mm) is typically 60 mm.

The H₂/HD spectra are monitored by 1+1' REMPI using the fourth harmonic of an unseeded Spectra Physics Nd:YAG laser (GCR-4, 266 nm, approx. 10 mJ per pulse) for the ionization step, and by applying a pulsed electric field (triggered after the laser pulse so that

Table 1. Error budget of the overall uncertainty in the calibration procedure.

Source	σ^a (MHz)
Doppler uncertainty ^b	60–100
AC Stark shift	30–50
Chirp shift	100
I ₂ calibration	10
Étalon drift	2
Statistics	80
Resulting uncertainty ^c	145–170

^a σ stands for standard deviation.

^b Estimated Doppler uncertainty at 91 000 cm⁻¹ for pure H₂.

^c Calculated as the square root of the sum of all specified uncertainties above.

laser excitation is performed under field-free conditions) to accelerate the ions towards an electron multiplier (Thorn EMI) detector. The distance between the interaction region and the ion detector is used as a time-of-flight mass spectrometer (perpendicular to both the VUV beam and the probe gas beam). The ion signal is recorded using a boxcar integrator (Stanford Research) to measure the area under a specific m/z peak in the time-of-flight mass spectrum. The signal is then transferred to a measuring program on a personal computer.

Possible effects of Doppler-induced shifts of the resonance frequency as a result of non-perpendicular alignment of the VUV and molecular beams is assessed by comparing spectral measurements obtained from pure hydrogen and from hydrogen/xenon mixtures, for which the molecular beam velocities are different. Furthermore it is noted that the interaction volume of the overlapping VUV/UV laser beams and the molecular beam is rather large, with the light beams possessing a diameter of 8 mm. Hence the AC-Stark induced shift of the ionizing UV beam is kept to a minimum. Finally there is the effect of frequency chirp in the dye amplifiers, which may result in a net shift between the centre frequency of the VUV beam used in the experiment, and the fundamental frequency (multiplied by a factor of 6) of the ring-dye laser, used for frequency calibration. All three issues are extensively discussed in the previous papers [24, 25] and similar procedures are followed to estimate resulting uncertainties associated with these systematic effects.

The VUV frequency is calibrated, using a fraction of the ring dye laser output to measure a Doppler-free absorption spectrum of molecular iodine and another fraction to record simultaneously an étalon transmission spectrum. The étalon is actively locked to a transmission mode of a frequency stabilized HeNe laser (Newport, model NL-1). The stability of the HeNe laser in the locked mode was measured with a fs frequency comb [30] and determined to be better than 2 MHz within a 20 min laser scan. The reference data for calibration were derived from a calculation of the iodine hyperfine structure by the IodineSpec program² An error budget of the calibration procedure is listed in table 1.

3. Results

The spectroscopic results are summarized in table 2 for the B ¹Σ_u⁺ ($v' = 0-2$) ← X ¹Σ_g⁺ ($v'' = 0$) Lyman bands in H₂ and in table 4 for the B ¹Σ⁺ ($v' = 0-2$) ← X ¹Σ⁺ ($v'' = 0$)

² The reference data for calibration were derived from a calculation of the iodine hyperfine structure by the IodineSpec program; see [31].

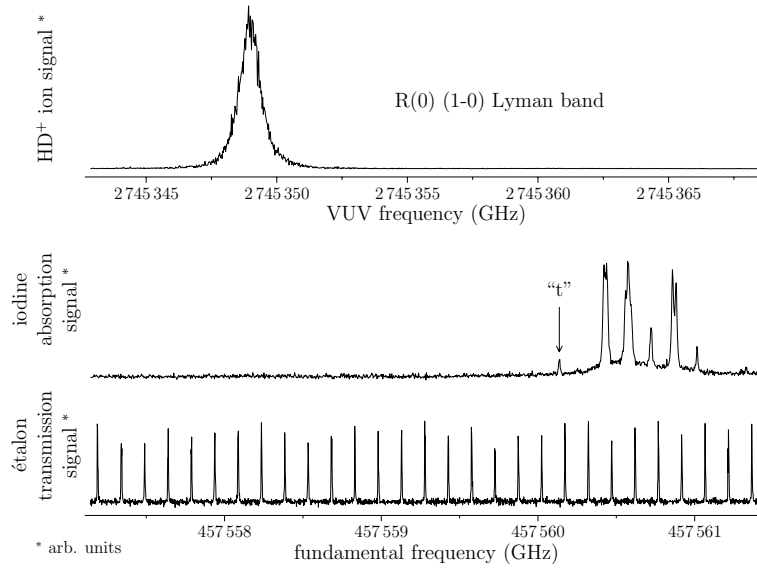


Figure 1. A spectrum of the $B^1\Sigma^+(v' = 1) \leftarrow X^1\Sigma^+(v'' = 0)$ R(0) Lyman transition in HD (upper trace). The linewidth of the measured transition is 900(25) MHz with a significant Doppler contribution. The lower two traces represent a simultaneously recorded Doppler-free calibration spectrum in molecular iodine (using the R(42) (5–5) transition with the ‘t’-component at 15 262.563 29 cm^{-1}) and an étalon transmission spectrum of a frequency stabilized étalon (lower trace).

Table 2. Transition wavenumbers of the Lyman bands in H_2 for $v' = 0\text{--}2$ with uncertainties given in parentheses (1σ). The values in the last three columns are in comparison with the reference values.

$\text{H}_2 B^1\Sigma_u^+(v', J') \leftarrow \text{H}_2 X^1\Sigma_g^+(v'' = 0, J'')$					
v'		$\frac{E_{\text{exp}}}{hc}$ (cm^{-1}) This work	$\frac{E_{\text{exp}} - E_{\text{ref}}}{hc}$ (cm^{-1})		
			[19] ^a	[18] ^b	[20] ^c
0	P(1)	90 085.0166 (50)		−0.0534	−0.0035
	R(0)	90 242.3373 (57)		0.0273	−0.0067
	R(1)	90 201.1545 (57)		−0.0255	−0.0046
1	P(1)	91 403.3354 (50)	0.1154	−0.0146	−0.0107
	P(2)	91 204.3491 (53)	1.1291	−0.0209	−0.0155
	P(3)	90 926.6825 (57)	0.1525	−0.0475	−0.0102
	R(0)	91 558.7222 (53)	−0.0878	0.0022	−0.0158
	R(1)	91 513.7115 (48)	0.0215	−0.0585	−0.0136
	R(2)	91 387.2454 (53)	−0.2146	−0.0346	−0.0352
2	R(3)	91 180.5323 (50)	0.0723	−0.0377	−0.0084
	R(2)	92 659.0480 (53)	−0.1320	0.0080	
	R(3)	92 446.1132 (57)	0.0432	−0.0368	
	R(4)	92 153.5373 (53)	−0.2127	−0.0127	

^a Uncertainties in [19] are specified as $\pm 0.003 \text{ \AA}$ at about 1100 \AA , which correspond to an uncertainty of approx. 0.25 cm^{-1} .

^b Calculated using the level energies given in [18] (table VI). The ground-state levels are taken from [15] (table 5). The specified uncertainties in [18] are less than 0.1 cm^{-1} .

^c Calculated using the level energies given in [20] (table I). The ground-state levels are calculated using the values of [32] (table 2, column B). The specified uncertainties in [20] are in the range of $\pm 0.005 \text{ cm}^{-1}$.

Table 3. Comparison of the difference frequencies of $R(J'') - P(J'' + 2)$ for $J'' = 0$ and $J'' = 1$ of the Lyman transitions in H_2 for $(v', v'') = (1, 0)$ (see table 2) and the $S(J'')$ transitions in the $X^1\Sigma_g^+$ ground state of H_2 (see [32]).

J''	$R(J'') - P(J'' + 2)$ (cm^{-1})	$S(J'')$ (cm^{-1})
	This work	[32]
0	354.373 (7)	354.3734 (2)
1	587.029 (7)	587.0325 (2)

Table 4. Transition wavenumbers of the Lyman bands in HD for $v' = 0-2$ with uncertainties given in parentheses (1σ). The reference data are taken from Dabrowski and Herzberg [16] using the values of table 10 for the $B^1\Sigma^+$ excited state of HD and from Evenson *et al* [33] and Ulivi *et al* [34] for the $X^1\Sigma^+$ ground state.

HD $B^1\Sigma^+(v', J') \leftarrow HD X^1\Sigma^+(v'' = 0, J'')$			
v'		$\frac{E}{hc}$ (cm^{-1})	$\frac{E_{exp} - E_{ref}}{hc}$ (cm^{-1})
0	P(1)	90 310.3785 (52)	-0.2536
	R(0)	90 428.9498 (46)	-0.2482
	R(1)	90 398.1855 (44)	-0.2445
	R(2)	90 307.5060 (44)	-0.2324
1	P(1)	91 457.7087 (44)	-0.1534
	P(2)	91 307.9139 (46)	-0.1545
	P(3)	91 098.5843 (45)	-0.1476
	R(0)	91 574.9856 (48)	-0.1525
	R(1)	91 541.6650 (45)	-0.1498
	R(2)	91 447.2193 (44)	-0.1319
2	R(3)	91 292.3011 (44)	-0.0861
	P(2)	92 425.8509 (44)	-0.2501
	R(0)	92 692.9203 (52)	-0.2503
	R(1)	92 657.4022 (52)	-0.2659

Lyman bands in HD. The upper trace in figure 1 shows the spectrum of the the $B^1\Sigma^+(v' = 1, J' = 1) \leftarrow X^1\Sigma^+(v'' = 0, J'' = 0)$ transition in HD. The lower two traces represent the calibration spectra consisting of a Doppler-free absorption spectrum of molecular iodine (the B-X system) and a transmission spectrum of an étalon.

In table 3, the differences $R(J'') - P(J'' + 2)$ for $J'' = 0$ and $J'' = 1$ of the $B^1\Sigma_u^+(v' = 1) \leftarrow X^1\Sigma_g^+(v'' = 0)$ Lyman transitions are compared with the corresponding $S(J'')$ transitions in the $X^1\Sigma_g^+$ ground state of H_2 . The values are consistent with the literature values of Jennings *et al* [32, 35]. This result provides an independent validation of the calibration procedures followed in the present study.

In the previous study by Philip *et al* [25], a value of $92\,659.092(4)\,cm^{-1}$ was reported for the R(2) line of the B-X (2,0) Lyman band, largely disagreeing with the present value. We note here that in [25] the very weak R(2) line was only detected in an overview scan and that its uncertainty was largely underestimated. The present value, listed in table 2, has a properly estimated uncertainty and should be regarded as the true value.

4. Conclusions

We present newly measured data of the Lyman transition of H_2 and HD in the lowest vibrational states of the $B^1\Sigma_u^+$ state ($v' = 0-2$) including an accurate calibration using Doppler-free

absorption spectroscopy. These measurements complete the recently published data set on the Lyman frequency transitions of Philip *et al* [25].

Comparisons of the experimental transition wavenumbers with reference data are listed in table 2 for H₂ and in table 4 for HD. In the case of H₂ we obtain an average deviation of (A) 0.089 cm⁻¹, (B) -0.026 cm⁻¹ and (C) -0.012 cm⁻¹ with respect to the reference data of (A) Baig and Connerade [19], measured by synchrotron spectroscopy, (B) Abgrall *et al* [18] and (C) Jungen *et al* [20], deduced from level energies. For HD, we compared our data with the values of Dabrowski and Herzberg [16] and obtain an average deviation of -0.19 cm⁻¹.

The present values on the transition frequencies are substantially more accurate than the previously published data and may find application in a number of astrophysical studies. The values are of high importance in particular for the comparison with transition frequencies of hydrogen lines observed in high-redshifted quasi-stellar objects.

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