

# VU Research Portal

## Evidence for $X(3872) \rightarrow \psi(2S)\gamma$ in $B\rightarrow K\gamma$ Decays and a Study of $B \rightarrow cc\gamma$

Aubert, B.; Zhang, L.L.; Bailey, D.S.; Raven, G.; Snoek, H.; Wang, W.; Liu, H.

### **published in**

Physical Review Letters

2009

### **DOI (link to publisher)**

[10.1103/PhysRevLett.102.132001](https://doi.org/10.1103/PhysRevLett.102.132001)

### **document version**

Publisher's PDF, also known as Version of record

[Link to publication in VU Research Portal](#)

### **citation for published version (APA)**

Aubert, B., Zhang, L. L., Bailey, D. S., Raven, G., Snoek, H., Wang, W., & Liu, H. (2009). Evidence for  $X(3872) \rightarrow \psi(2S)\gamma$  in  $B\rightarrow K\gamma$  Decays and a Study of  $B \rightarrow cc\gamma$ . *Physical Review Letters*, 102(13), 1-7. Article 132001. <https://doi.org/10.1103/PhysRevLett.102.132001>

### **General rights**

Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

- Users may download and print one copy of any publication from the public portal for the purpose of private study or research.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying the publication in the public portal

### **Take down policy**

If you believe that this document breaches copyright please contact us providing details, and we will remove access to the work immediately and investigate your claim.

### **E-mail address:**

[vuresearchportal.ub@vu.nl](mailto:vuresearchportal.ub@vu.nl)

## Evidence for $X(3872) \rightarrow \psi(2S)\gamma$ in $B^\pm \rightarrow X(3872)K^\pm$ Decays and a Study of $B \rightarrow c\bar{c}\gamma K$

- B. Aubert,<sup>1</sup> M. Bona,<sup>1</sup> Y. Karyotakis,<sup>1</sup> J. P. Lees,<sup>1</sup> V. Poireau,<sup>1</sup> E. Prencipe,<sup>1</sup> X. Prudent,<sup>1</sup> V. Tisserand,<sup>1</sup> J. Garra Tico,<sup>2</sup> E. Grauges,<sup>2</sup> L. Lopez,<sup>3a,3b</sup> A. Palano,<sup>3a,3b</sup> M. Pappagallo,<sup>3a,3b</sup> G. Eigen,<sup>4</sup> B. Stugu,<sup>4</sup> L. Sun,<sup>4</sup> G. S. Abrams,<sup>5</sup> M. Battaglia,<sup>5</sup> D. N. Brown,<sup>5</sup> R. N. Cahn,<sup>5</sup> R. G. Jacobsen,<sup>5</sup> L. T. Kerth,<sup>5</sup> Yu. G. Kolomensky,<sup>5</sup> G. Lynch,<sup>5</sup> I. L. Osipenkov,<sup>5</sup> M. T. Ronan,<sup>5,\*</sup> K. Tackmann,<sup>5</sup> T. Tanabe,<sup>5</sup> C. M. Hawkes,<sup>6</sup> N. Soni,<sup>6</sup> A. T. Watson,<sup>6</sup> H. Koch,<sup>7</sup> T. Schroeder,<sup>7</sup> D. Walker,<sup>8</sup> D. J. Asgeirsson,<sup>9</sup> B. G. Fulsom,<sup>9</sup> C. Hearty,<sup>9</sup> T. S. Mattison,<sup>9</sup> J. A. McKenna,<sup>9</sup> M. Barrett,<sup>10</sup> A. Khan,<sup>10</sup> V. E. Blinov,<sup>11</sup> A. D. Bukan,<sup>11</sup> A. R. Buzykaev,<sup>11</sup> V. P. Druzhinin,<sup>11</sup> V. B. Golubev,<sup>11</sup> A. P. Onuchin,<sup>11</sup> S. I. Serednyakov,<sup>11</sup> Yu. I. Skovpen,<sup>11</sup> E. P. Solodov,<sup>11</sup> K. Yu. Todyshev,<sup>11</sup> M. Bondioli,<sup>12</sup> S. Curry,<sup>12</sup> I. Eschrich,<sup>12</sup> D. Kirkby,<sup>12</sup> A. J. Lankford,<sup>12</sup> P. Lund,<sup>12</sup> M. Mandelkern,<sup>12</sup> E. C. Martin,<sup>12</sup> D. P. Stoker,<sup>12</sup> S. Abachi,<sup>13</sup> C. Buchanan,<sup>13</sup> J. W. Gary,<sup>14</sup> F. Liu,<sup>14</sup> O. Long,<sup>14</sup> B. C. Shen,<sup>14,\*</sup> G. M. Vitug,<sup>14</sup> Z. Yasin,<sup>14</sup> L. Zhang,<sup>14</sup> V. Sharma,<sup>15</sup> C. Campagnari,<sup>16</sup> T. M. Hong,<sup>16</sup> D. Kovalskyi,<sup>16</sup> M. A. Mazur,<sup>16</sup> J. D. Richman,<sup>16</sup> T. W. Beck,<sup>17</sup> A. M. Eisner,<sup>17</sup> C. J. Flacco,<sup>17</sup> C. A. Heusch,<sup>17</sup> J. Kroseberg,<sup>17</sup> W. S. Lockman,<sup>17</sup> A. J. Martinez,<sup>17</sup> T. Schalk,<sup>17</sup> B. A. Schumm,<sup>17</sup> A. Seiden,<sup>17</sup> M. G. Wilson,<sup>17</sup> L. O. Winstrom,<sup>17</sup> C. H. Cheng,<sup>18</sup> D. A. Doll,<sup>18</sup> B. Echenard,<sup>18</sup> F. Fang,<sup>18</sup> D. G. Hitlin,<sup>18</sup> I. Narsky,<sup>18</sup> T. Piatenko,<sup>18</sup> F. C. Porter,<sup>18</sup> R. Andreassen,<sup>19</sup> G. Mancinelli,<sup>19</sup> B. T. Meadows,<sup>19</sup> K. Mishra,<sup>19</sup> M. D. Sokoloff,<sup>19</sup> P. C. Bloom,<sup>20</sup> W. T. Ford,<sup>20</sup> A. Gaz,<sup>20</sup> J. F. Hirschauer,<sup>20</sup> M. Nagel,<sup>20</sup> U. Nauenberg,<sup>20</sup> J. G. Smith,<sup>20</sup> K. A. Ulmer,<sup>20</sup> S. R. Wagner,<sup>20</sup> R. Ayad,<sup>21,\*</sup> A. Soffer,<sup>21</sup> W. H. Toki,<sup>21</sup> R. J. Wilson,<sup>21</sup> D. D. Altenburg,<sup>22</sup> E. Feltresi,<sup>22</sup> A. Hauke,<sup>22</sup> H. Jasper,<sup>22</sup> M. Karbach,<sup>22</sup> J. Merkel,<sup>22</sup> A. Petzold,<sup>22</sup> B. Spaan,<sup>22</sup> K. Wacker,<sup>22</sup> M. J. Kobel,<sup>23</sup> W. F. Mader,<sup>23</sup> R. Nogowski,<sup>23</sup> K. R. Schubert,<sup>23</sup> R. Schwierz,<sup>23</sup> A. Volk,<sup>23</sup> D. Bernard,<sup>24</sup> G. R. Bonneau,<sup>24</sup> E. Latour,<sup>24</sup> M. Verderi,<sup>24</sup> P. J. Clark,<sup>25</sup> S. Playfer,<sup>25</sup> J. E. Watson,<sup>25</sup> M. Andreotti,<sup>26a,26b</sup> D. Bettom,<sup>26a</sup> C. Bozzi,<sup>26a</sup> R. Calabrese,<sup>26a,26b</sup> A. Cecchi,<sup>26a,26b</sup> G. Cibinetto,<sup>26a,26b</sup> P. Franchini,<sup>26a,26b</sup> E. Luppi,<sup>26a,26b</sup> M. Negrini,<sup>26a,26b</sup> A. Petrella,<sup>26a,26b</sup> L. Piemontese,<sup>26a</sup> V. Santoro,<sup>26a,26b</sup> R. Baldini-Ferroli,<sup>27</sup> A. Calcaterra,<sup>27</sup> R. de Sangro,<sup>27</sup> G. Finocchiaro,<sup>27</sup> S. Pacetti,<sup>27</sup> P. Patteri,<sup>27</sup> I. M. Peruzzi,<sup>27,\*</sup> M. Piccolo,<sup>27</sup> M. Rama,<sup>27</sup> A. Zallo,<sup>27</sup> A. Buzzo,<sup>28a</sup> R. Contri,<sup>28a,28b</sup> M. Lo Vetere,<sup>28a,28b</sup> M. M. Macri,<sup>28a</sup> M. R. Monge,<sup>28a,28b</sup> S. Passaggio,<sup>28a</sup> C. Patrignani,<sup>28a,28b</sup> E. Robutti,<sup>28a</sup> A. Santroni,<sup>28a,28b</sup> S. Tosi,<sup>28a,28b</sup> K. S. Chaisanguanthum,<sup>29</sup> M. Morii,<sup>29</sup> A. Adametz,<sup>30</sup> J. Marks,<sup>30</sup> S. Schenk,<sup>30</sup> U. Uwer,<sup>30</sup> V. Klose,<sup>31</sup> H. M. Lacker,<sup>31</sup> D. J. Bard,<sup>32</sup> P. D. Dauncey,<sup>32</sup> J. A. Nash,<sup>32</sup> M. Tibbetts,<sup>32</sup> P. K. Behera,<sup>33</sup> X. Chai,<sup>33</sup> M. J. Charles,<sup>33</sup> U. Mallik,<sup>33</sup> J. Cochran,<sup>34</sup> H. B. Crowley,<sup>34</sup> L. Dong,<sup>34</sup> W. T. Meyer,<sup>34</sup> S. Prell,<sup>34</sup> E. I. Rosenberg,<sup>34</sup> A. E. Rubin,<sup>34</sup> Y. Y. Gao,<sup>35</sup> A. V. Gritsan,<sup>35</sup> Z. J. Guo,<sup>35</sup> C. K. Lae,<sup>35</sup> N. Arnaud,<sup>36</sup> J. Béquilleux,<sup>36</sup> A. D'Orazio,<sup>36</sup> M. Davier,<sup>36</sup> J. Firmino da Costa,<sup>36</sup> G. Grosdidier,<sup>36</sup> A. Höcker,<sup>36</sup> V. Lepeltier,<sup>36</sup> F. Le Diberder,<sup>36</sup> A. M. Lutz,<sup>36</sup> S. Pruvot,<sup>36</sup> P. Roudeau,<sup>36</sup> M. H. Schune,<sup>36</sup> J. Serrano,<sup>36</sup> V. Sordini,<sup>36,\*</sup> A. Stocchi,<sup>36</sup> G. Wormser,<sup>36</sup> D. J. Lange,<sup>37</sup> D. M. Wright,<sup>37</sup> I. Bingham,<sup>38</sup> J. P. Burke,<sup>38</sup> C. A. Chavez,<sup>38</sup> J. R. Fry,<sup>38</sup> E. Gabathuler,<sup>38</sup> R. Gamet,<sup>38</sup> D. E. Hutchcroft,<sup>38</sup> D. J. Payne,<sup>38</sup> C. Touramanis,<sup>38</sup> A. J. Bevan,<sup>39</sup> C. K. Clarke,<sup>39</sup> K. A. George,<sup>39</sup> F. Di Lodovico,<sup>39</sup> R. Sacco,<sup>39</sup> M. Sigamani,<sup>39</sup> G. Cowan,<sup>40</sup> H. U. Flaecher,<sup>40</sup> D. A. Hopkins,<sup>40</sup> S. Paramesvaran,<sup>40</sup> F. Salvatore,<sup>40</sup> A. C. Wren,<sup>40</sup> D. N. Brown,<sup>41</sup> C. L. Davis,<sup>41</sup> A. G. Denig,<sup>42</sup> M. Fritsch,<sup>42</sup> W. Gradl,<sup>42</sup> G. Schott,<sup>42</sup> K. E. Alwyn,<sup>43</sup> D. Bailey,<sup>43</sup> R. J. Barlow,<sup>43</sup> Y. M. Chia,<sup>43</sup> C. L. Edgar,<sup>43</sup> G. Jackson,<sup>43</sup> G. D. Lafferty,<sup>43</sup> T. J. West,<sup>43</sup> J. I. Yi,<sup>43</sup> J. Anderson,<sup>44</sup> C. Chen,<sup>44</sup> A. Jawahery,<sup>44</sup> D. A. Roberts,<sup>44</sup> G. Simi,<sup>44</sup> J. M. Tuggle,<sup>44</sup> C. Dallapiccola,<sup>45</sup> X. Li,<sup>45</sup> E. Salvati,<sup>45</sup> S. Saremi,<sup>45</sup> R. Cowan,<sup>46</sup> D. Dujmic,<sup>46</sup> P. H. Fisher,<sup>46</sup> G. Sciolla,<sup>46</sup> M. Spitznagel,<sup>46</sup> F. Taylor,<sup>46</sup> R. K. Yamamoto,<sup>46</sup> M. Zhao,<sup>46</sup> P. M. Patel,<sup>47</sup> S. H. Robertson,<sup>47</sup> A. Lazzaro,<sup>48a,48b</sup> V. Lombardo,<sup>48a</sup> F. Palombo,<sup>48a,48b</sup> J. M. Bauer,<sup>49</sup> L. Cremaldi,<sup>49</sup> R. Godang,<sup>49,\*</sup> R. Kroeger,<sup>49</sup> D. A. Sanders,<sup>49</sup> D. J. Summers,<sup>49</sup> H. W. Zhao,<sup>49</sup> M. Simard,<sup>50</sup> P. Taras,<sup>50</sup> F. B. Viaud,<sup>50</sup> H. Nicholson,<sup>51</sup> G. De Nardo,<sup>52a,52b</sup> L. Lista,<sup>52a</sup> D. Monorchio,<sup>52a,52b</sup> G. Onorato,<sup>52a,52b</sup> C. Sciacca,<sup>52a,52b</sup> G. Raven,<sup>53</sup> H. L. Snoek,<sup>53</sup> C. P. Jessop,<sup>54</sup> K. J. Knoepfel,<sup>54</sup> J. M. Lo Secco,<sup>54</sup> W. F. Wang,<sup>54</sup> G. Benelli,<sup>55</sup> L. A. Corwin,<sup>55</sup> K. Honscheid,<sup>55</sup> H. Kagan,<sup>56</sup> R. Kass,<sup>55</sup> J. P. Morris,<sup>55</sup> A. M. Rahimi,<sup>55</sup> J. J. Regensburger,<sup>55</sup> S. J. Sekula,<sup>55</sup> Q. K. Wong,<sup>55</sup> N. L. Blount,<sup>56</sup> J. Brau,<sup>56</sup> R. Frey,<sup>56</sup> O. Igonkina,<sup>56</sup> J. A. Kolb,<sup>56</sup> M. Lu,<sup>56</sup> R. Rahmat,<sup>56</sup> N. B. Sinev,<sup>56</sup> D. Strom,<sup>56</sup> J. Strube,<sup>56</sup> E. Torrence,<sup>56</sup> G. Castelli,<sup>57a,57b</sup> N. Gagliardi,<sup>57a,57b</sup> M. Margoni,<sup>57a,57b</sup> M. Morandin,<sup>57a</sup> M. Posocco,<sup>57a</sup> M. Rotondo,<sup>57a</sup> F. Simonetto,<sup>57a,57b</sup> R. Stroili,<sup>57a,57b</sup> C. Voci,<sup>57a,57b</sup> P. del Amo Sanchez,<sup>58</sup> E. Ben-Haim,<sup>58</sup> H. Briand,<sup>58</sup> G. Calderini,<sup>58</sup> J. Chauveau,<sup>58</sup> P. David,<sup>58</sup> L. Del Buono,<sup>58</sup> O. Hamon,<sup>58</sup> Ph. Leruste,<sup>58</sup> J. Ocariz,<sup>58</sup> A. Perez,<sup>58</sup> J. Prendki,<sup>58</sup> S. Sitt,<sup>58</sup> L. Gladney,<sup>59</sup> M. Biasini,<sup>60a,60b</sup> R. Covarelli,<sup>60a,60b</sup> E. Manoni,<sup>60a,60b</sup> C. Angelini,<sup>61a,61b</sup> G. Batignani,<sup>61a,61b</sup> S. Bettarini,<sup>61a,61b</sup> M. Carpinelli,<sup>61a,61b,\*</sup> A. Cervelli,<sup>61a,61b</sup> F. Forti,<sup>61a,61b</sup> M. A. Giorgi,<sup>61a,61b</sup> A. Lusiani,<sup>61a,61c</sup> G. Marchiori,<sup>61a,61b</sup> M. Morganti,<sup>61a,61b</sup> N. Neri,<sup>61a,61b</sup> E. Paoloni,<sup>61a,61b</sup> G. Rizzo,<sup>61a,61b</sup> J. J. Walsh,<sup>61a,61b</sup> D. Lopes Pegna,<sup>62</sup> C. Lu,<sup>62</sup> J. Olsen,<sup>62</sup> A. J. S. Smith,<sup>62</sup> A. V. Telnov,<sup>62</sup> F. Anulli,<sup>63a</sup> E. Baracchini,<sup>63a,63b</sup> G. Cavoto,<sup>63a</sup> D. del Re,<sup>63a,63b</sup> E. Di Marco,<sup>63a,63b</sup> R. Faccini,<sup>63a,63b</sup>

- F. Ferrarotto,<sup>63a</sup> F. Ferroni,<sup>63a,63b</sup> M. Gaspero,<sup>63a,63b</sup> P. D. Jackson,<sup>63a</sup> L. Li Gioi,<sup>63a</sup> M. A. Mazzoni,<sup>63a</sup> S. Morganti,<sup>63a</sup> G. Piredda,<sup>63a</sup> F. Polci,<sup>63a,63b</sup> F. Renga,<sup>63a,63b</sup> C. Voena,<sup>63a</sup> M. Ebert,<sup>64</sup> T. Hartmann,<sup>64</sup> H. Schröder,<sup>64</sup> R. Waldi,<sup>64</sup> T. Adye,<sup>65</sup> B. Franek,<sup>65</sup> E. O. Olaiya,<sup>65</sup> F. F. Wilson,<sup>65</sup> S. Emery,<sup>66</sup> M. Escalier,<sup>66</sup> L. Esteve,<sup>66</sup> S. F. Ganzhur,<sup>66</sup> G. Hamel de Monchenault,<sup>66</sup> W. Kozanecki,<sup>66</sup> G. Vasseur,<sup>66</sup> Ch. Yèche,<sup>66</sup> M. Zito,<sup>66</sup> X. R. Chen,<sup>67</sup> H. Liu,<sup>67</sup> W. Park,<sup>67</sup> M. V. Purohit,<sup>67</sup> R. M. White,<sup>67</sup> J. R. Wilson,<sup>67</sup> M. T. Allen,<sup>68</sup> D. Aston,<sup>68</sup> R. Bartoldus,<sup>68</sup> P. Bechtle,<sup>68</sup> J. F. Benitez,<sup>68</sup> R. Cenci,<sup>68</sup> J. P. Coleman,<sup>68</sup> M. R. Convery,<sup>68</sup> J. C. Dingfelder,<sup>68</sup> J. Dorfan,<sup>68</sup> G. P. Dubois-Felsmann,<sup>68</sup> W. Dunwoodie,<sup>68</sup> R. C. Field,<sup>68</sup> A. M. Gabareen,<sup>68</sup> S. J. Gowdy,<sup>68</sup> M. T. Graham,<sup>68</sup> P. Grenier,<sup>68</sup> C. Hast,<sup>68</sup> W. R. Innes,<sup>68</sup> J. Kaminski,<sup>68</sup> M. H. Kelsey,<sup>68</sup> H. Kim,<sup>68</sup> P. Kim,<sup>68</sup> M. L. Kocian,<sup>68</sup> D. W. G. S. Leith,<sup>68</sup> S. Li,<sup>68</sup> B. Lindquist,<sup>68</sup> S. Luitz,<sup>68</sup> V. Luth,<sup>68</sup> H. L. Lynch,<sup>68</sup> D. B. MacFarlane,<sup>68</sup> H. Marsiske,<sup>68</sup> R. Messner,<sup>68</sup> D. R. Muller,<sup>68</sup> H. Neal,<sup>68</sup> S. Nelson,<sup>68</sup> C. P. O'Grady,<sup>68</sup> I. Ofte,<sup>68</sup> A. Perazzo,<sup>68</sup> M. Perl,<sup>68</sup> B. N. Ratcliff,<sup>68</sup> A. Roodman,<sup>68</sup> A. A. Salnikov,<sup>68</sup> R. H. Schindler,<sup>68</sup> J. Schwiening,<sup>68</sup> A. Snyder,<sup>68</sup> D. Su,<sup>68</sup> M. K. Sullivan,<sup>68</sup> K. Suzuki,<sup>68</sup> S. K. Swain,<sup>68</sup> J. M. Thompson,<sup>68</sup> J. Va'vra,<sup>68</sup> A. P. Wagner,<sup>68</sup> M. Weaver,<sup>68</sup> C. A. West,<sup>68</sup> W. J. Wisniewski,<sup>68</sup> M. Wittgen,<sup>68</sup> D. H. Wright,<sup>68</sup> H. W. Wulsin,<sup>68</sup> A. K. Yarritu,<sup>68</sup> K. Yi,<sup>68</sup> C. C. Young,<sup>68</sup> V. Ziegler,<sup>68</sup> P. R. Burchat,<sup>69</sup> A. J. Edwards,<sup>69</sup> S. A. Majewski,<sup>69</sup> T. S. Miyashita,<sup>69</sup> B. A. Petersen,<sup>69</sup> L. Wilden,<sup>69</sup> S. Ahmed,<sup>70</sup> M. S. Alam,<sup>70</sup> J. A. Ernst,<sup>70</sup> B. Pan,<sup>70</sup> M. A. Saeed,<sup>70</sup> S. B. Zain,<sup>70</sup> S. M. Spanier,<sup>71</sup> B. J. Wogsland,<sup>71</sup> R. Eckmann,<sup>72</sup> J. L. Ritchie,<sup>72</sup> A. M. Ruland,<sup>72</sup> C. J. Schilling,<sup>72</sup> R. F. Schwitters,<sup>72</sup> B. W. Drummond,<sup>73</sup> J. M. Izen,<sup>73</sup> X. C. Lou,<sup>73</sup> F. Bianchi,<sup>74a,74b</sup> D. Gamba,<sup>74a,74b</sup> M. Pelliccioni,<sup>74a,74b</sup> M. Bomben,<sup>75a,75b</sup> L. Bosisio,<sup>75a,75b</sup> C. Cartaro,<sup>75a,75b</sup> G. Della Ricca,<sup>75a,75b</sup> L. Lanceri,<sup>75a,75b</sup> L. Vitale,<sup>75a,75b</sup> V. Azzolini,<sup>76</sup> N. Lopez-March,<sup>76</sup> F. Martinez-Vidal,<sup>76</sup> D. A. Milanes,<sup>76</sup> A. Oyanguren,<sup>76</sup> J. Albert,<sup>77</sup> Sw. Banerjee,<sup>77</sup> B. Bhuyan,<sup>77</sup> H. H. F. Choi,<sup>77</sup> K. Hamano,<sup>77</sup> R. Kowalewski,<sup>77</sup> M. J. Lewczuk,<sup>77</sup> I. M. Nugent,<sup>77</sup> J. M. Roney,<sup>77</sup> R. J. Sobie,<sup>77</sup> T. J. Gershon,<sup>78</sup> P. F. Harrison,<sup>78</sup> J. Ilic,<sup>78</sup> T. E. Latham,<sup>78</sup> G. B. Mohanty,<sup>78</sup> H. R. Band,<sup>79</sup> X. Chen,<sup>79</sup> S. Dasu,<sup>79</sup> K. T. Flood,<sup>79</sup> Y. Pan,<sup>79</sup> M. Pierini,<sup>79</sup> R. Prepost,<sup>79</sup> C. O. Vuosalo,<sup>79</sup> and S. L. Wu<sup>79</sup>

(The BABAR Collaboration)

<sup>1</sup>Laboratoire de Physique des Particules, IN2P3/CNRS et Université de Savoie, F-74941 Annecy-Le-Vieux, France<sup>2</sup>Universitat de Barcelona, Facultat de Fisica, Departament ECM, E-08028 Barcelona, Spain<sup>3a</sup>INFN Sezione di Bari, I-70126 Bari, Italy<sup>3b</sup>Dipartimento di Fisica, Università di Bari, I-70126 Bari, Italy<sup>4</sup>University of Bergen, Institute of Physics, N-5007 Bergen, Norway<sup>5</sup>Lawrence Berkeley National Laboratory and University of California, Berkeley, California 94720, USA<sup>6</sup>University of Birmingham, Birmingham, B15 2TT, United Kingdom<sup>7</sup>Ruhr Universität Bochum, Institut für Experimentalphysik 1, D-44780 Bochum, Germany<sup>8</sup>University of Bristol, Bristol BS8 1TL, United Kingdom<sup>9</sup>University of British Columbia, Vancouver, British Columbia, Canada V6T 1Z1<sup>10</sup>Brunel University, Uxbridge, Middlesex UB8 3PH, United Kingdom<sup>11</sup>Budker Institute of Nuclear Physics, Novosibirsk 630090, Russia<sup>12</sup>University of California at Irvine, Irvine, California 92697, USA<sup>13</sup>University of California at Los Angeles, Los Angeles, California 90024, USA<sup>14</sup>University of California at Riverside, Riverside, California 92521, USA<sup>15</sup>University of California at San Diego, La Jolla, California 92093, USA<sup>16</sup>University of California at Santa Barbara, Santa Barbara, California 93106, USA<sup>17</sup>University of California at Santa Cruz, Institute for Particle Physics, Santa Cruz, California 95064, USA<sup>18</sup>California Institute of Technology, Pasadena, California 91125, USA<sup>19</sup>University of Cincinnati, Cincinnati, Ohio 45221, USA<sup>20</sup>University of Colorado, Boulder, Colorado 80309, USA<sup>21</sup>Colorado State University, Fort Collins, Colorado 80523, USA<sup>22</sup>Technische Universität Dortmund, Fakultät Physik, D-44221 Dortmund, Germany<sup>23</sup>Technische Universität Dresden, Institut für Kern- und Teilchenphysik, D-01062 Dresden, Germany<sup>24</sup>Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, F-91128 Palaiseau, France<sup>25</sup>University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom<sup>26a</sup>INFN Sezione di Ferrara, I-44100 Ferrara, Italy<sup>26b</sup>Dipartimento di Fisica, Università di Ferrara, I-44100 Ferrara, Italy<sup>27</sup>INFN Laboratori Nazionali di Frascati, I-00044 Frascati, Italy<sup>28a</sup>INFN Sezione di Genova, I-16146 Genova, Italy<sup>28b</sup>Dipartimento di Fisica, Università di Genova, I-16146 Genova, Italy<sup>29</sup>Harvard University, Cambridge, Massachusetts 02138, USA

<sup>30</sup>*Universität Heidelberg, Physikalisches Institut, Philosophenweg 12, D-69120 Heidelberg, Germany*<sup>31</sup>*Humboldt-Universität zu Berlin, Institut für Physik, Newtonstr. 15, D-12489 Berlin, Germany*<sup>32</sup>*Imperial College London, London, SW7 2AZ, United Kingdom*<sup>33</sup>*University of Iowa, Iowa City, Iowa 52242, USA*<sup>34</sup>*Iowa State University, Ames, Iowa 50011-3160, USA*<sup>35</sup>*Johns Hopkins University, Baltimore, Maryland 21218, USA*<sup>36</sup>*Laboratoire de l'Accélérateur Linéaire, IN2P3/CNRS et Université Paris-Sud 11, Centre Scientifique d'Orsay, B. P. 34, F-91898 Orsay Cedex, France*<sup>37</sup>*Lawrence Livermore National Laboratory, Livermore, California 94550, USA*<sup>38</sup>*University of Liverpool, Liverpool L69 7ZE, United Kingdom*<sup>39</sup>*Queen Mary, University of London, London, E1 4NS, United Kingdom*<sup>40</sup>*University of London, Royal Holloway and Bedford New College, Egham, Surrey TW20 0EX, United Kingdom*<sup>41</sup>*University of Louisville, Louisville, Kentucky 40292, USA*<sup>42</sup>*Johannes Gutenberg-Universität Mainz, Institut für Kernphysik, D-55099 Mainz, Germany*<sup>43</sup>*University of Manchester, Manchester M13 9PL, United Kingdom*<sup>44</sup>*University of Maryland, College Park, Maryland 20742, USA*<sup>45</sup>*University of Massachusetts, Amherst, Massachusetts 01003, USA*<sup>46</sup>*Massachusetts Institute of Technology, Laboratory for Nuclear Science, Cambridge, Massachusetts 02139, USA*<sup>47</sup>*McGill University, Montréal, Québec, Canada H3A 2T8*<sup>48a</sup>*INFN Sezione di Milano, I-20133 Milano, Italy*<sup>48b</sup>*Dipartimento di Fisica, Università di Milano, I-20133 Milano, Italy*<sup>49</sup>*University of Mississippi, University, Mississippi 38677, USA*<sup>50</sup>*Université de Montréal, Physique des Particules, Montréal, Québec, Canada H3C 3J7*<sup>51</sup>*Mount Holyoke College, South Hadley, Massachusetts 01075, USA*<sup>52a</sup>*INFN Sezione di Napoli, I-80126 Napoli, Italy*<sup>52b</sup>*Dipartimento di Scienze Fisiche, Università di Napoli Federico II, I-80126 Napoli, Italy*<sup>53</sup>*NIKHEF, National Institute for Nuclear Physics and High Energy Physics, NL-1009 DB Amsterdam, The Netherlands*<sup>54</sup>*University of Notre Dame, Notre Dame, Indiana 46556, USA*<sup>55</sup>*Ohio State University, Columbus, Ohio 43210, USA*<sup>56</sup>*University of Oregon, Eugene, Oregon 97403, USA*<sup>57a</sup>*INFN Sezione di Padova, I-35131 Padova, Italy*<sup>57b</sup>*Dipartimento di Fisica, Università di Padova, I-35131 Padova, Italy*<sup>58</sup>*Laboratoire de Physique Nucléaire et de Hautes Energies, IN2P3/CNRS, Université Pierre et Marie Curie-Paris6, Université Denis Diderot-Paris7, F-75252 Paris, France*<sup>59</sup>*University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*<sup>60a</sup>*INFN Sezione di Perugia, I-06100 Perugia, Italy*<sup>60b</sup>*Dipartimento di Fisica, Università di Perugia, I-06100 Perugia, Italy*<sup>61a</sup>*INFN Sezione di Pisa, I-56127 Pisa, Italy*<sup>61b</sup>*Dipartimento di Fisica, Università di Pisa, I-56127 Pisa, Italy*<sup>61c</sup>*Scuola Normale Superiore di Pisa, I-56127 Pisa, Italy*<sup>62</sup>*Princeton University, Princeton, New Jersey 08544, USA*<sup>63a</sup>*INFN Sezione di Roma, I-00185 Roma, Italy*<sup>63b</sup>*Dipartimento di Fisica, Università di Roma La Sapienza, I-00185 Roma, Italy*<sup>64</sup>*Universität Rostock, D-18051 Rostock, Germany*<sup>65</sup>*Rutherford Appleton Laboratory, Chilton, Didcot, Oxon, OX11 0QX, United Kingdom*<sup>66</sup>*CEA, Ifu, SPP, Centre de Saclay, F-91191 Gif-sur-Yvette, France*<sup>67</sup>*University of South Carolina, Columbia, South Carolina 29208, USA*<sup>68</sup>*Stanford Linear Accelerator Center, Stanford, California 94309, USA*<sup>69</sup>*Stanford University, Stanford, California 94305-4060, USA*<sup>70</sup>*State University of New York, Albany, New York 12222, USA*<sup>71</sup>*University of Tennessee, Knoxville, Tennessee 37996, USA*<sup>72</sup>*University of Texas at Austin, Austin, Texas 78712, USA*<sup>73</sup>*University of Texas at Dallas, Richardson, Texas 75083, USA*<sup>74a</sup>*INFN Sezione di Torino, I-10125 Torino, Italy*<sup>74b</sup>*Dipartimento di Fisica Sperimentale, Università di Torino, I-10125 Torino, Italy*<sup>75a</sup>*INFN Sezione di Trieste, I-34127 Trieste, Italy*<sup>75b</sup>*Dipartimento di Fisica, Università di Trieste, I-34127 Trieste, Italy*<sup>76</sup>*IFIC, Universitat de Valencia-CSIC, E-46071 Valencia, Spain*<sup>77</sup>*University of Victoria, Victoria, British Columbia, Canada V8W 3P6*<sup>78</sup>*Department of Physics, University of Warwick, Coventry CV4 7AL, United Kingdom*<sup>79</sup>*University of Wisconsin, Madison, Wisconsin 53706, USA*

(Received 2 September 2008; published 30 March 2009)

In a search for  $B \rightarrow c\bar{c}\gamma K$  decays with the *BABAR* detector, where  $c\bar{c}$  includes  $J/\psi$  and  $\psi(2S)$ , and  $K$  includes  $K^\pm$ ,  $K_S^0$ , and  $K^*(892)$ , we find evidence for  $X(3872) \rightarrow J/\psi\gamma$  and  $X(3872) \rightarrow \psi(2S)\gamma$  with  $3.6\sigma$  and  $3.5\sigma$  significance, respectively. We measure the product of branching fractions  $\mathcal{B}(B^\pm \rightarrow X(3872)K^\pm) \times \mathcal{B}(X(3872) \rightarrow J/\psi\gamma) = [2.8 \pm 0.8(\text{stat}) \pm 0.1(\text{syst})] \times 10^{-6}$  and  $\mathcal{B}(B^\pm \rightarrow X(3872)K^\pm) \times \mathcal{B}(X(3872) \rightarrow \psi(2S)\gamma) = [9.5 \pm 2.7(\text{stat}) \pm 0.6(\text{syst})] \times 10^{-6}$ .

DOI: 10.1103/PhysRevLett.102.132001

PACS numbers: 13.20.Gd, 13.20.He, 14.40.Gx

The  $X(3872)$  state discovered by the Belle Collaboration in the decay  $B^\pm \rightarrow K^\pm X(3872)$ ,  $X(3872) \rightarrow J/\psi\pi^+\pi^-$  [1] is now well established [2]. *BABAR* has seen evidence for the decay  $X(3872) \rightarrow J/\psi\gamma$ , which implies positive  $C$ -parity [3]. A variety of theoretical interpretations [4] exist for this state, including conventional charmonium interpretations [5] and exotic QCD proposals such as a  $\bar{D}^0 D^{*0}$  molecule [6]. While  $\bar{D}^0 D^{*0}$  molecular proposals can accommodate decays to  $J/\psi\gamma$ , the branching fraction for decays to  $\psi(2S)\gamma$  is expected to be very small [7]. These models allow for the possibility of an admixture of a  $\bar{D}^0 D^{*0}$  bound state with, for example, a  $c\bar{c}$  meson. Because the  $\chi_{c1}(2P)$  state potentially decays to  $\psi(2S)\gamma$  at a rate many times higher than to  $J/\psi\gamma$ , the decay  $X(3872) \rightarrow \psi(2S)\gamma$  could be enhanced due to  $c\bar{c}-\bar{D}^0 D^{*0}$  mixing.

We present a study of the decay  $B \rightarrow XK$ , where the notation  $X$  represents any state decaying radiatively to  $J/\psi\gamma$  or  $\psi(2S)\gamma$  [the  $\chi_{c1,2}$  and  $X(3872)$  states in particular], and  $K$  encompasses  $K^\pm$ ,  $K_S^0$ ,  $K^{*\pm}(892)$ , and  $K^{*0}(892)$ . We consider  $J/\psi$  mesons decaying to  $e^+e^-$  or  $\mu^+\mu^-$ , and  $\psi(2S)$  decaying to  $e^+e^-$ ,  $\mu^+\mu^-$ , or  $J/\psi\pi^+\pi^-$ . Kaons are required to decay to final states consisting of charged particles:  $K_S^0 \rightarrow \pi^+\pi^-$ ,  $K^{*\pm} \rightarrow K_S^0(\pi^+\pi^-)\pi^\pm$ , and  $K^{*0} \rightarrow K^\pm\pi^\mp$ .

The data sample for this analysis consists of  $(465 \pm 5)$  million  $B\bar{B}$  pairs collected with the *BABAR* detector at the PEP-II asymmetric  $e^+e^-$  collider at SLAC. This represents  $424 \text{ fb}^{-1}$  of data taken at the  $Y(4S)$  resonance. The *BABAR* detector is described in detail elsewhere [8]. The event selection, determined independently from the data, is based on Monte Carlo (MC) simulated events with the aim of maximizing significance.

The  $J/\psi$  candidates are formed using pairs of leptons whose invariant mass is in the range  $(2.96, 3.15) \text{ GeV}/c^2$  for electrons (including bremsstrahlung photons) and  $(3.06, 3.13) \text{ GeV}/c^2$  for muons. For  $\psi(2S) \rightarrow \ell^+\ell^-$ , the candidate invariant masses are required to be in the range  $(3.61, 3.73) \text{ GeV}/c^2$  for electrons or  $(3.65, 3.72) \text{ GeV}/c^2$  for muons. The  $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$  candidates are composed of  $J/\psi$  candidates decaying as described but with a tighter mass requirement of  $(3.01, 3.15) \text{ GeV}/c^2$  for the  $e^+e^-$  decay mode. To form a  $\psi(2S)$  candidate, the  $J/\psi$  candidate is mass constrained to the nominal PDG value [9] and combined with a pair of oppositely charged tracks requiring  $(0.4, 0.6) \text{ GeV}/c^2$  and  $(3.68, 3.69) \text{ GeV}/c^2$  for the dipion and  $\psi(2S)$  invariant masses, respectively. All

four final decay particles are constrained to the same decay vertex. Electrons are identified by a likelihood-based selector with  $>92\%$  efficiency and negligible fake rate. Muons are selected by a neural net process with  $>85\%$  efficiency and a  $\pi(K)$  fake rate of  $<6\%(<10\%)$ . Pions are drawn from the list of all charged tracks in the event.

We reconstruct  $X \rightarrow c\bar{c}\gamma$  candidates from a mass-constrained  $J/\psi$  ( $\psi(2S)$ ) candidate combined with a photon with an energy greater than  $30(100) \text{ MeV}$ . Additional selection criteria are applied to the shape of the lateral distribution ( $0.001 < \text{LAT} < 0.5$ ) [10] and azimuthal asymmetry (as measured by the Zernike moment  $A_{42} < 0.1$ ) [11] of the photon-shower energy. For  $X \rightarrow J/\psi\gamma$ , the radiative  $\gamma$  candidate is rejected if, when combined with any other  $\gamma$  from the event, it has an invariant mass consistent with the  $\pi^0$  mass,  $124 < m_{\gamma\gamma} < 146 \text{ MeV}/c^2$ .

The  $K_S^0$  candidates are required to be within  $\pm 17 \text{ MeV}/c^2$  of the nominal  $K_S^0$  mass [9], and the significance of the distance of the reconstructed decay vertex from the primary vertex must be greater than  $3.7$  standard deviations ( $\sigma$ ). The excited kaons are required to have an invariant mass within the range  $0.7 < m(K^*) < 1.1 \text{ GeV}/c^2$ . For  $K_S^0$ ,  $K^{*\pm}$ , and  $K^{*0}$  candidates associated with  $X \rightarrow \psi(2S)\gamma$ , additional requirements are placed on the  $\chi^2$  vertex probability of the kaon,  $P(\chi^2) > 0.001, 0.02$ , and  $0.002$ , respectively. Kaons are chosen by a likelihood-based selector with an efficiency of  $\sim 95\%$  and misidentification rates of  $\sim 5\%$ ,  $\sim 4\%$ , and  $< 10\%$  for  $\pi, \mu$ , and  $p$ , respectively, over the momentum range in this analysis.

We form the final  $B$  candidate from an  $X$  candidate and a kaon constrained to originate from the same vertex. To identify  $B$  candidates, we use two kinematic variables,  $m_B$  and  $m_{\text{miss}}$ . The unconstrained mass of the reconstructed  $B$  candidate is  $m_B = \sqrt{E_B^2/c^4 - p_B^2/c^2}$ , where  $E_B$  and  $p_B$  are obtained by summing the energies and momenta of the particles in the candidate  $B$  meson. The missing mass is defined as  $m_{\text{miss}} = \sqrt{(p_{e^+e^-} - \hat{p}_B)^2/c^2}$ , where  $p_{e^+e^-}$  is the four-momentum of the beam  $e^+e^-$  system and  $\hat{p}_B$  is the four-momentum of the  $B$  candidate after applying a  $B$  mass constraint. For  $X \rightarrow J/\psi[\psi(2S)]\gamma$  events, we require  $m_B$  to be within  $^{+30}_{-36}(\pm 20) \text{ MeV}/c^2$  of the nominal  $B$  mass [9]. Our  $B$  candidate selection is further refined by imposing criteria on the  $\chi^2$  probability for the  $B$  vertex: for all  $X \rightarrow J/\psi\gamma$  modes,  $P(\chi^2) > 0.0001$ , and for  $X \rightarrow \psi(2S)\gamma$  modes,  $P(\chi^2) > 0.01, 0.002$ , and  $0.05$  for the  $K^\pm, K_S^0$ ,

and  $K^*$  modes, respectively. The ratio of the second and zeroth Fox-Wolfram moments ( $R_2 < 0.45$ ) [12] is used to separate isotropic  $B$  events from continuum background events. Once a  $B$  candidate has been established, it and its daughter decays are refit with the  $B$  mass constrained to the known value [9].

We perform a one(two)-dimensional unbinned extended maximum-likelihood (UML) fit to  $m_{\text{miss}}$  (and  $m_{K^*}$ , if applicable), and use the  $s$ Plot formalism [13] to project our signal events into  $m_X$ , the invariant mass of the  $X$  candidate. This is a background-subtraction technique by weighting each event based on how signal- or background-like it is. The  $s$ Plot displays the number of  $B \rightarrow XK$  signal-like events as a function of  $m_X$ . We extract the signal yield for a given decay mode by fitting this resultant  $m_X$  distribution with shapes for signal and background determined from MC simulation.

The signal event probability density functions (PDFs) are determined from MC-simulated  $B \rightarrow \chi_{c1}K$  and  $B \rightarrow X(3872)K$  events. Only reconstructed events exactly matching the generated decay chain particles are used to parameterize the signal PDFs. The PDF shapes for  $B \rightarrow \chi_{c2}K$  are the same as for  $\chi_{c1}$ , with the below-noted exception of the  $m_X$  distribution. The  $m_{\text{miss}}$  distribution is modeled with a Crystal Ball function [14],  $m_X$  with a single Gaussian for the  $\chi_{c2}$  decay modes and narrower core Gaussian plus a second wider Gaussian sharing the same

mean for all other signal modes, and  $m_{K^*}$  with the convolution of a Breit-Wigner and a Gaussian.

The background PDFs are determined from fits to generic  $B^+B^-$ ,  $B^0\bar{B}^0$ ,  $q\bar{q}$ , and  $\tau^+\tau^-$  MC samples, and are dominated by events from  $B\bar{B}$  decays that include a  $J/\psi$  or  $\psi(2S)$  in their decay chain. For the  $B^\pm \rightarrow XK^\pm$  and  $B^0 \rightarrow XK_S^0$  decay modes, the background in  $m_{\text{miss}}$  consists of two parts: a nonpeaking combinatoric component modeled with an ARGUS function [15], and a peaking component that shares the Crystal Ball parameterization used for signal events. These backgrounds are modeled as linear in  $m_X$ . The  $K^*$  decay modes have three background components: events that peak in  $m_{\text{miss}}$  but are flat in  $m_{K^*}$  (“nonresonant”) and vice versa (“ $K^*$  combinatoric”), and those that do not peak in either distribution (“combinatoric”). The peaking  $m_{\text{miss}}$  and  $m_{K^*}$  distributions use the same parameterization and values found by fitting to the signal MC sample. The nonpeaking  $m_{\text{miss}}$  distributions are fit with an ARGUS function, while the nonpeaking  $m_{K^*}$  distribution is modeled with a linear function. Both combinatoric background types are flat in  $m_X$ , while the nonresonant backgrounds (typically  $B \rightarrow XK\pi$ ) have a flat and peaking component in  $m_X$ . However, because none of these background events are signal-like in both  $m_{\text{miss}}$  and  $m_{K^*}$ , they are not present in the  $s$ Plot projection in  $m_X$ .

To account for potential differences between data and MC calculations, the values for the  $m_{\text{miss}}$  ARGUS and  $m_X$

TABLE I. Summary of the analysis results.  $N_S$  is the bias-corrected number of signal events extracted from the  $m_X$   $s$ Plot,  $\sigma$  is the total significance of the signal yield  $N_S$  measured in standard deviations (statistical and systematic uncertainties combined in quadrature) from the null result,  $\epsilon$  is the total efficiency for the decay mode, and derived  $\mathcal{B}$  is the measurement (with 90% confidence level upper limit [16]) of  $\mathcal{B}(B \rightarrow \chi_{c1,2}K)$  or  $\mathcal{B}(B \rightarrow X(3872)K) \times \mathcal{B}(X \rightarrow c\bar{c}\gamma)$ . Uncertainties are statistical and systematic, respectively.

Decay	$N_S$	$\sigma$	$\epsilon(\%)$	Derived $\mathcal{B}$
$\chi_{c1}$				$\times 10^{-4}$
$\chi_{c1}K^\pm$	$1018 \pm 34 \pm 14$	28	11.0	$4.5 \pm 0.1 \pm 0.3$
$\chi_{c1}K^0$	$242 \pm 16 \pm 5$	14	8.7	$4.2 \pm 0.3 \pm 0.3$
$\chi_{c1}K^{*\pm}$	$71 \pm 13 \pm 8$	4.7	5.7	$2.6 \pm 0.5 \pm 0.4$
$\chi_{c1}K^{*0}$	$255 \pm 25 \pm 11$	9.5	7.9	$2.5 \pm 0.2 \pm 0.2$
$\chi_{c2}$				$\times 10^{-5}$
$\chi_{c2}K^\pm$	$14.0 \pm 7.9 \pm 1.1$	1.8	12.3	$1.0 \pm 0.6 \pm 0.1 (< 1.8)$
$\chi_{c2}K^0$	$6.1 \pm 3.9 \pm 1.1$	1.5	11.1	$1.5 \pm 0.9 \pm 0.3 (< 2.8)$
$\chi_{c2}K^{*\pm}$	$1.2 \pm 4.7 \pm 6.1$	0.2	4.2	$1.1 \pm 4.3 \pm 5.5 (< 12)$
$\chi_{c2}K^{*0}$	$38.8 \pm 10.5 \pm 1.1$	3.7	8.3	$6.6 \pm 1.8 \pm 0.5$
$X(3872)(J/\psi\gamma)$				$\times 10^{-6}$
$XK^\pm$	$23.0 \pm 6.4 \pm 0.6$	3.6	14.5	$2.8 \pm 0.8 \pm 0.1$
$XK^0$	$5.3 \pm 3.6 \pm 0.2$	1.5	11.0	$2.6 \pm 1.8 \pm 0.2 (< 4.9)$
$XK^{*\pm}$	$0.6 \pm 2.3 \pm 0.1$	0.3	6.9	$0.7 \pm 2.6 \pm 0.1 (< 4.8)$
$XK^{*0}$	$2.8 \pm 5.2 \pm 0.4$	0.5	10.4	$0.7 \pm 1.4 \pm 0.1 (< 2.8)$
$X(3872)(\psi(2S)\gamma)$				$\times 10^{-6}$
$XK^\pm$	$25.4 \pm 7.3 \pm 0.7$	3.5	10.4	$9.5 \pm 2.7 \pm 0.6$
$XK^0$	$8.0 \pm 3.9 \pm 0.5$	2.0	8.4	$11.4 \pm 5.5 \pm 1.0 (< 19)$
$XK^{*\pm}$	$1.9 \pm 2.9 \pm 2.9$	0.5	5.0	$6.4 \pm 9.8 \pm 9.6 (< 28)$
$XK^{*0}$	$-1.4 \pm 3.3 \pm 0.3$	...	6.7	$-1.3 \pm 3.1 \pm 0.3 (< 4.4)$

linear parameters for the background events are left as free parameters in the final fit to data. We also allow the height of the  $m_X$  Gaussian peaks to float, which we use to derive the number of signal events.

The effectiveness of the signal extraction method is validated on fully simulated MC events for  $\chi_{c1,2}$  and  $X(3872)$  signal events, with random samples generated from the MC background distribution. Successful performance of the fit is verified on simulated datasets assuming the number of signal and background events from the known branching fractions and efficiencies. We apply small corrections ( $<5\%$ ) to account for bias in the results of the MC fit validation.

We determine the efficiency from the fraction of the events generated in MC simulation that survive the analysis selection criteria and are returned by the fitting procedure. We calculate the branching fraction for each decay mode using  $\mathcal{B}(B \rightarrow XK) = N_S / (N_{B\bar{B}} \times \epsilon \times f)$  where  $N_S$  is the bias-corrected number of signal events from the fit to the  $m_X$ ,  $N_{B\bar{B}}$  is the number of  $B\bar{B}$  pairs in the data set,  $\epsilon$  is the total signal extraction efficiency, and  $f$  represents all secondary branching fractions. The fit results, efficiencies, and derived branching fractions are summarized in Table I.

For most of the  $B \rightarrow X(3872)K$  decay channels, the largest source of systematic uncertainty affecting the signal yield comes from the uncertainty in the true  $X(3872)$  mass and width ( $\sim 2\%$  for the  $K^\pm$  modes). In the case of the  $K^\pm$  and  $K_S^0$  decay modes for  $X(3872) \rightarrow \psi(2S)\gamma$ , an alternate parametrization of the  $m_X$  shape was considered for background events, as indicated by the MC simulation. A correction equal to half the difference between the results of the two background model choices, with a systematic error equal to this amount, is applied to the final result. This is the largest yield-related systematic uncertainty for the  $X(3872)(\psi(2S)\gamma)K^\pm$  mode ( $\sim 2\%$ ). For  $B \rightarrow \chi_{c1,2}K$ , uncertainty in the fit bias, PDF parameters, and MC-data differences for the mean value  $m_X$  for signal events all contribute in varying though roughly equal amounts.

Regarding systematic uncertainties related to the branching fraction calculations, one of the main contributors is the total uncertainty associated with the identification of all particle types ( $\sim 4\%$ ). The uncertainty in secondary branching fractions, beyond the control of this analysis, is the dominant systematic uncertainty for  $\mathcal{B}(B \rightarrow \chi_{c1}K)$  and  $\mathcal{B}(B^0 \rightarrow \chi_{c2}K^{*0})$  ( $\sim 6\%$ ). Effects from tracking, photon corrections and  $B$  counting are also considered, but are all less than 2%.

Figure 1 shows the fit to  $m_X$  in the mass range  $3.411 < m_X < 3.611 \text{ GeV}/c^2$ . We observe all of the expected  $B \rightarrow \chi_{c1}K$  decay modes, in good agreement with previous measurements. We find  $3.7\sigma$  evidence for  $B^0 \rightarrow \chi_{c2}K^{*0}$ , and set upper limits for the remaining  $B \rightarrow \chi_{c2}K$  decays.

Fits to  $m_X$  in the range  $3.772 < m_X < 3.972 \text{ GeV}/c^2$  are shown in Fig. 2 for decays to  $J/\psi\gamma$ . We confirm evidence

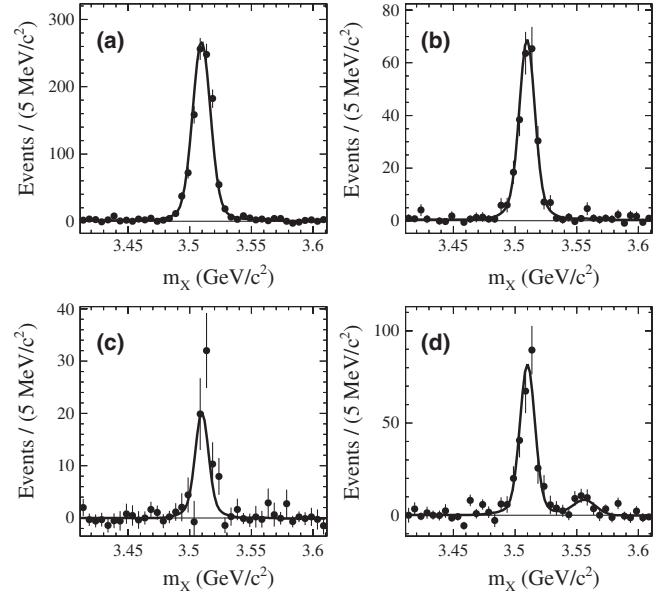


FIG. 1. Plot of the number of signal events versus  $m_X$  for (a)  $B^\pm \rightarrow \chi_{c1,2} K^\pm$ , (b)  $B^0 \rightarrow \chi_{c1,2} K_S^0$ , (c)  $B^\pm \rightarrow \chi_{c1,2} K^{*\pm}$ , and (d)  $B^0 \rightarrow \chi_{c1,2} K^{*0}$ . The solid curve is the fit to the data.

for the decay  $X(3872) \rightarrow J/\psi\gamma$  in  $B^\pm \rightarrow X(3872)K^\pm$ , measuring  $\mathcal{B}(B^\pm \rightarrow X(3872)K^\pm) \times \mathcal{B}(X(3872) \rightarrow J/\psi\gamma) = [2.8 \pm 0.8(\text{stat}) \pm 0.1(\text{syst})] \times 10^{-6}$  with a significance of  $3.6\sigma$ . This value is in good agreement with the previous *BABAR* result [3], which it supersedes, and represents the most precise measurement of this branching

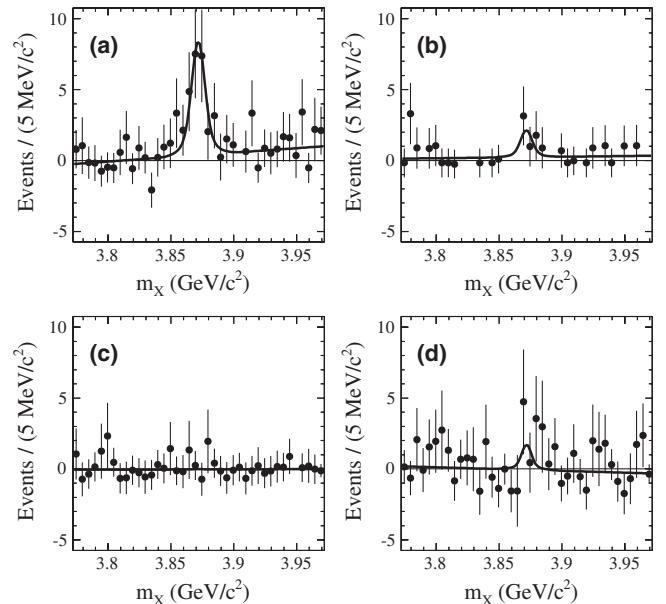


FIG. 2. Plot of the number of extracted signal events versus  $m_X$  for (a)  $B^\pm \rightarrow X(3872)K^\pm$ , (b)  $B^0 \rightarrow X(3872)K_S^0$ , (c)  $B^\pm \rightarrow X(3872)K^{*\pm}$ , and (d)  $B^0 \rightarrow X(3872)K^{*0}$ , where  $X(3872) \rightarrow J/\psi\gamma$ . The solid curve is the fit to the data.

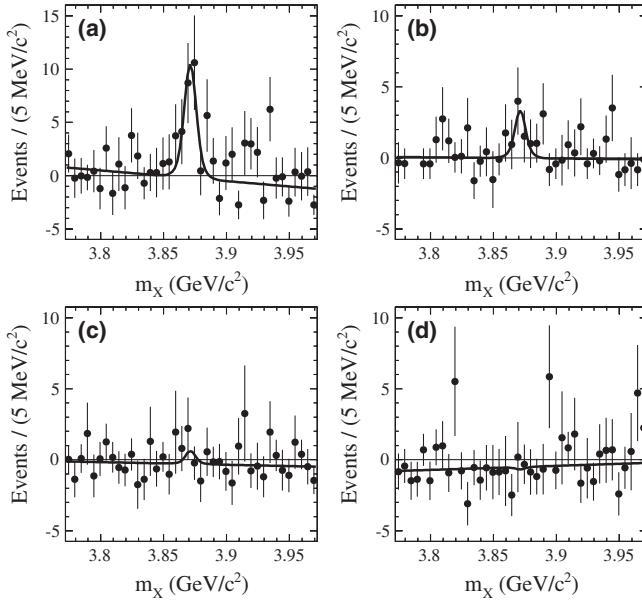


FIG. 3. Plot of the number of extracted signal events versus  $m_X$  for (a)  $B^\pm \rightarrow X(3872)K^\pm$ , (b)  $B^0 \rightarrow X(3872)K_S^0$ , (c)  $B^\pm \rightarrow X(3872)K^{*\pm}$ , and (d)  $B^0 \rightarrow X(3872)K^{*0}$ , where  $X(3872) \rightarrow \psi(2S)\gamma$ . The solid curve is the fit to the data.

fraction to date. We find no significant signal in the other decay modes.

Figure 3 shows the fit to the  $m_X$  distribution for  $X \rightarrow \psi(2S)\gamma$  in the range  $3.772 < m_X < 3.972 \text{ GeV}/c^2$ . In our search for  $X(3872) \rightarrow \psi(2S)\gamma$  in  $B^\pm \rightarrow X(3872)K^\pm$ , we find the first evidence for this decay with a significance of  $3.5\sigma$ . We derive  $\mathcal{B}(B^\pm \rightarrow X(3872)K^\pm) \times \mathcal{B}(X(3872) \rightarrow \psi(2S)\gamma) = [9.5 \pm 2.7(\text{stat}) \pm 0.6(\text{syst})] \times 10^{-6}$ . We find no significant signals in the other decay modes.

To search for other new resonances, the  $m_X$  invariant mass window is extended up to the kinematic limit. Even after combining all decay modes together, there are no indications of any further signals above the prominent  $X(3872)$  peak.

In summary, we present first evidence for the decay  $X(3872) \rightarrow \psi(2S)\gamma$ , and updated measurements of the  $X(3872) \rightarrow J/\psi\gamma$  and  $B \rightarrow \chi_{c1,2}K$  decays. We find evidence for the factorization-suppressed [17] decay  $B^0 \rightarrow \chi_{c2}K^{*0}$ , but see no evidence for other  $\chi_{c2}K$  decays. Taking the statistical and systematic errors in quadrature, we find a ratio of  $\frac{\mathcal{B}(X(3872) \rightarrow \psi(2S)\gamma)}{\mathcal{B}(X(3872) \rightarrow J/\psi\gamma)} = 3.4 \pm 1.4$ . This relatively large branching fraction for  $X(3872) \rightarrow \psi(2S)\gamma$  is generally inconsistent with a purely  $\bar{D}^0 D^{*0}$  molecular interpretation of the  $X(3872)$ , and possibly indicates mixing with a significant  $c\bar{c}$  component.

We are grateful for the excellent luminosity and machine conditions provided by our PEP-II colleagues, and for the substantial dedicated effort from the computing organizations that support *BABAR*. The collaborating institutions wish to thank SLAC for its support and kind hospitality.

This work is supported by DOE and NSF (USA), NSERC (Canada), CEA and CNRS-IN2P3 (France), BMBF and DFG (Germany), INFN (Italy), FOM (The Netherlands), NFR (Norway), MES (Russia), MEC (Spain), and STFC (United Kingdom). Individuals have received support from the Marie Curie EIF (European Union) and the A. P. Sloan Foundation.

\*Deceased

<sup>†</sup>Now at Temple University, Philadelphia, PA 19122, USA

<sup>‡</sup>Now at Tel Aviv University, Tel Aviv, 69978, Israel

<sup>§</sup>Also with Università di Perugia, Dipartimento di Fisica, Perugia, Italy

<sup>||</sup>Also with Università di Roma La Sapienza, I-00185 Roma, Italy

<sup>¶</sup>Now at University of South Alabama, Mobile, AL 36688, USA

<sup>\*\*</sup>Also with Università di Sassari, Sassari, Italy

- [1] S. K. Choi *et al.* (Belle Collaboration), Phys. Rev. Lett. **91**, 262001 (2003).
- [2] D. Acosta *et al.* (CDF Collaboration), Phys. Rev. Lett. **93**, 072001 (2004); V. M. Abazov *et al.* (D0 Collaboration), *ibid.* **93**, 162002 (2004); B. Aubert *et al.* (*BABAR* Collaboration), Phys. Rev. D **71**, 071103 (2005).
- [3] B. Aubert *et al.* (*BABAR* Collaboration), Phys. Rev. D **74**, 071101(R) (2006).
- [4] S. Godfrey and S. Olsen, Annu. Rev. Nucl. Part. Sci. **58**, 51 (2008).
- [5] T. Barnes and S. Godfrey, Phys. Rev. D **69**, 054008 (2004).
- [6] N. A. Tornqvist, Phys. Lett. B **590**, 209 (2004); E. S. Swanson, *ibid.* **598**, 197 (2004); M. B. Voloshin, *ibid.* **579**, 316 (2004); E. Braaten and M. Kusunoki, Phys. Rev. D **71**, 074005 (2005).
- [7] E. S. Swanson, Phys. Rep. **429**, 243 (2006).
- [8] B. Aubert *et al.* (*BABAR* Collaboration), Nucl. Instrum. Methods Phys. Res., Sect. A **479**, 1 (2002).
- [9] C. Amsler *et al.* (Particle Data Group), Phys. Lett. B **667**, 1 (2008).
- [10] A. Drescher *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **237**, 464 (1985).
- [11] R. Sinkus and T. Voss, Nucl. Instrum. Methods Phys. Res., Sect. A **391**, 360 (1997).
- [12] G. C. Fox and S. Wolfram, Nucl. Phys. B **149**, 413 (1979).
- [13] M. Pivk and F. R. Le Diberder, Nucl. Instrum. Methods Phys. Res., Sect. A **555**, 356 (2005).
- [14] J. E. Gaiser, Report No. SLAC-R-255, Appendix F, 1982.
- [15] H. Albrecht *et al.* (ARGUS Collaboration), Phys. Lett. B **185**, 218 (1987); **241**, 278 (1990).
- [16] The upper limit (UL) is calculated from  $\int_0^{UL} G(x)/\int_0^{+\infty} G(x) = 0.9$ , where  $G(x)$  is a Gaussian with mean equal to the central value of the branching fraction measurement and standard deviation equal to the total uncertainty.
- [17] M. Suzuki, Phys. Rev. D **66**, 037503 (2002); P. Colangelo, F. De Fazio, and T. N. Pham, Phys. Lett. B **542**, 71 (2002).